

# LIGHT QUARK MASSES FROM RBC & UKQCD COLLABORATIONS

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# Light Quark Mass Results

- \* arXiv/1011.0892: 2+1 flavor dynamical domain-wall fermions with continuum limit

$$\begin{aligned}m_{ud}^{\overline{MS}}(2GeV) &= 3.59(13)_{\text{stat}}(14)_{\text{syst}}(8)_{\text{ren}} \text{ MeV}, \\m_s^{\overline{MS}}(2GeV) &= 96.2(1.6)_{\text{stat}}(0.2)_{\text{syst}}(2.1)_{\text{ren}} \text{ MeV}\end{aligned}$$



- \* PRD78(2008)114509: 2+1 flavor dynamical domain-wall fermions w/o continuum lim

$$\begin{aligned}m_{ud}^{\overline{MS}}(2GeV) &= 3.72(13)_{\text{stat}}(18)_{\text{syst}}(33)_{\text{ren}} \text{ MeV}, \\m_s^{\overline{MS}}(2GeV) &= 107.3(4.4)_{\text{stat}}(4.9)_{\text{syst}}(9.7)_{\text{ren}} \text{ MeV}\end{aligned}$$

- \* How these improvements are achieved is reviewed....



# Plan

- \* results (now and then)
- \* quark mass improvement: step by step
  - \* 2+1 flavor dynamical domain-wall fermion simulation with 2 lattice spacing
  - \* non-perturbative renormalization with RI/SMOM schemes
- \* results
- \* outlook
  
- \* related papers
  - \* PRD78(2008)114509, arXiv/1011.0892 by RBC/UKQCD
  - \* arXiv/1006.0422 by R. Arthur and P. Boyle



# at Year 2008

- \* PRD78(2008)114509: 2+1 flavor dynamical domain-wall fermions w/o continuum lim
- \* Iwasaki gauge + DWF:  $\beta=2.13$  ( $a \approx 0.11$  fm),  $24^3 \times 64$  ( $L_s=2.7$  fm)
- \* Pros
  - \* chiral fermion (simpler handling of chiral extrapolation)
  - \* RI/MOM scheme non-perturbative renormalization (NPR)
- \* Cons
  - \* single lattice spacing
  - \* large NPR error
- \* arXiv/1011.0892: 2+1 flavor dynamical domain-wall fermions with continuum limit



# at Year 2010

- \* arXiv/1011.0892: 2+1 flavor dynamical domain-wall fermions with continuum limit
- \* Iwasaki gauge + DWF:  $\beta=2.13$  ( $a \approx 0.11$  fm),  $24^3 \times 64$  ( $L_s=2.7$  fm) and
- \* Iwasaki gauge + DWF:  $\beta=2.25$  ( $a \approx 0.9$  fm),  $32^3 \times 64$  ( $L_s=2.8$  fm)
- \* What's new:
  - \* 2 lattice spacings
    - \* continuum limit
    - \* more robust chiral extrapolation possible
  - \* RI/SMOM scheme NPR
    - \* much reduced systematic error



# 2+1 flavor dynamical chiral fermions

$m_s a$	$m_l a$	$\tilde{m}_s / \tilde{m}_l$	$\Delta t_{light}$	$\tau(\text{Ref.}[1])$	$\tau(\text{MD})$	Acceptance	$\langle P \rangle$	$\langle \bar{\psi} \psi(m_l) \rangle$
$V/a = 24^3 \times 64, L_s = 16, \beta = 2.13, a^{-1} = 1.73(3) \text{ GeV}, m_{res} a = 0.003152(43), \tau/\text{traj} = 1$								
0.04	0.005	5.3	1/6	4460	8980	73%	0.588053(4)	0.001224(2)
	0.01	3.3	1/5	5020	8540	70%	0.588009(5)	0.001738(2)
$V/a = 32^3 \times 64, L_s = 16, \beta = 2.25, a^{-1} = 2.28(3) \text{ GeV}, m_{res} a = 0.0006664(76), \tau/\text{traj} = 2$								
0.03	0.004	6.6	1/8	—	6856	72%	0.615587(3)	0.000673(1)
	0.006	4.6	1/8	—	7650	76%	0.615585(3)	0.000872(1)
	0.008	3.5	1/7	—	5930	73%	0.615571(4)	0.001066(1)

- \* Iwasaki gauge and domain-wall fermions
- \* 2 lattice spacings
- \* 1 strange mass & 2-3 u. d quark masses
- \* combined chiral and continuum extrapolation is performed



# extrapolation / interpolation

- \* strange mass:
  - \* valence: 2: one unitary and one smaller
  - \* sea: many points by reweighting
  - \* combining these makes 2nd unitary point
  - \* 2nd point is chosen so to interpolate
- \* u, d quark mass:
  - \* valence: 3-4 points to cover  $225 \text{ MeV} \leq m_\pi \leq 420 \text{ MeV}$
  - \* sea: 2-3 points to cover  $290 \text{ MeV} \leq m_\pi \leq 420 \text{ MeV}$
  - \* extrapolation to  $m_\pi \rightarrow 135 \text{ MeV}$
- \* lattice spacing:  $a=0.114 \text{ fm}$  &  $0.087 \text{ fm}$  to take the continuum limit



# determining the physical point

- \* need to determine ( $a$ ,  $m_{ud}$ ,  $m_s$ ): “physical point” for each lattice spacing
- \* using mass of  $\pi$ ,  $K$ ,  $\Omega$
- \*  $m_\pi$ ,  $m_K$ 
  - \* statistically most precisely calculated quantity
  - \* chiral behavior known better than others
- \*  $m_\Omega$ : mass of sss baryon
  - \* reasonably well controlled statistical error
  - \* Chiral extrapolation is easy (no chiral log)
- \* matching the continuum-extrapolated lattice results
  - \* physical point is determined
- \* now you can predict the other quantities



# chiral and continuum global fit

- \* mass renormalization needed to handle data with multi lattice spacing
- \* renormalization scale for each lattice depend on the lattice scale determined from the chiral fit
- \* decoupling these two simplifies the whole calculation framework
- \* fixed trajectory method



# notation

- \* **coarser** lattice is referred to as **32** cubed lattice
- \* **finer** lattice is referred to as **24** cubed lattice
- \*  $m_l$ : light quark mass for u, d
- \*  $m_h$ : heavy quark mass for s
- \*  $m_{ll}$ : “pion” with  $(m_l, m_l)$
- \*  $m_{lh}$ : “kaon” with  $(m_l, m_h)$
- \*  $m_{hhh}$ : “ $\Omega$ ” with  $(m_h, m_h, m_h)$



# primary lattice ensemble and other

- \* define the lattice scheme with the primary ensemble (lattice spacing)
- \* in this work 32 cubed is the primary ensemble
- \* do matching the other ensembles to primary
- \* in this work we match 24 cubed to 32 cubed
- \* all results are parametrized with the parameter in the primary ensemble
- \* all 24 and 32 cubed data are parametrized in 32 cubed parameter
- \* do chiral fit with  $a^2$  error taken into account
- \* input mass of  $\pi$ , K,  $\Omega \Rightarrow$  primary ensemble physical point is determined
- \* do renormalization of the primary ensemble parameter: to get quark mass
- \* (we do need renormalization for another ensemble to get rid of a-error)



# parameterization

- \* everything is parametrized in 32 cubed (finer) lattice scheme:
- \* introduce
  - \* ratio of lattice spacing  $R_a = a_{32}/a_{24}$
  - \* ratio of mass renormalization constant
    - \*  $Z_l = Z_m^{24}/Z_m^{32}$  for (u, d)
    - \*  $Z_h = Z_m^{24}/Z_m^{32}$  for (s)
    - \*  $Z_m$  to match to a mass independent scheme
    - \* difference arises due to mass dependent lattice artifact  $O(a^2)$
    - \* function of mass, but, in practice, one value has an “allowed range”
- \* determined iteratively at a simulated mass on the finer lattice
- \* then 24 cubed (coarser) lattice can be treated in the same ground as 32



# fixed trajectory scheme

- \* renormalized trajectory obtained by matching to a simulated point of the primary ensemble
- \* ex: use 32 cubed  $(m_l, m_h) = (0.006, 0.03)$

<b>M</b>	$(am_l)^M$	$(am_h)^M$	$(am_l)^e$	$(am_h)^e$	$Z_l$	$Z_h$	$R_a$
$32^3$	0.004	0.03	0.00313(13)	0.03812(80)	0.980(15)	0.976(11)	0.7617(72)
$32^3$	0.006	0.03	0.00583(12)	0.03839(51)	0.981(9)	0.974(7)	0.7583(46)
$32^3$	0.008	0.03	0.00860(19)	0.03869(64)	0.979(10)	0.972(8)	0.7545(58)
$24^3$	0.005	0.04	0.00545(11)	0.03148(51)	0.985(12)	0.978(9)	0.7620(57)
$24^3$	0.01	0.04	0.00897(18)	0.03074(57)	0.974(11)	0.968(9)	0.7517(70)

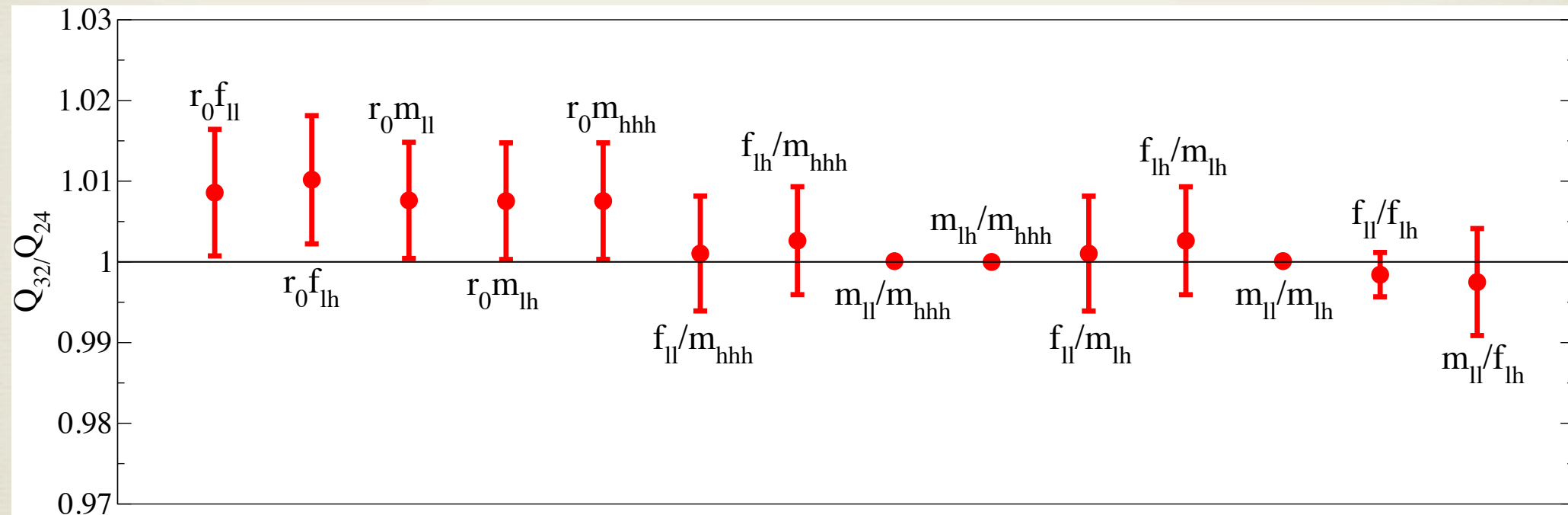
- \* stable over the range of simulated quark mass



# a scaling test at along the trajectory

- \*  $Q_{32}/Q_{24}$  for  $Q$  being a dimension less ratio of observables
- \*  $Q_{32}/Q_{24} = 1 \pm 0.01$  level (1 for perfect scaling)

matched with  $m_l=0.006$ ,  $m_h=0.03$  for 32 cubed





# global fit

- \* Now that we determined the necessary parameters, we can perform the global fit using  $24^3$  and  $32^3$
- \* 2 fit types
  - \* ChPT (NLO)
    - \* good if simulated masses are in the chiral regime
    - \* with and without finite volume correction to assess finite V error
  - \* analytic (1st order)
    - \* good if the chiral regime is below physical point

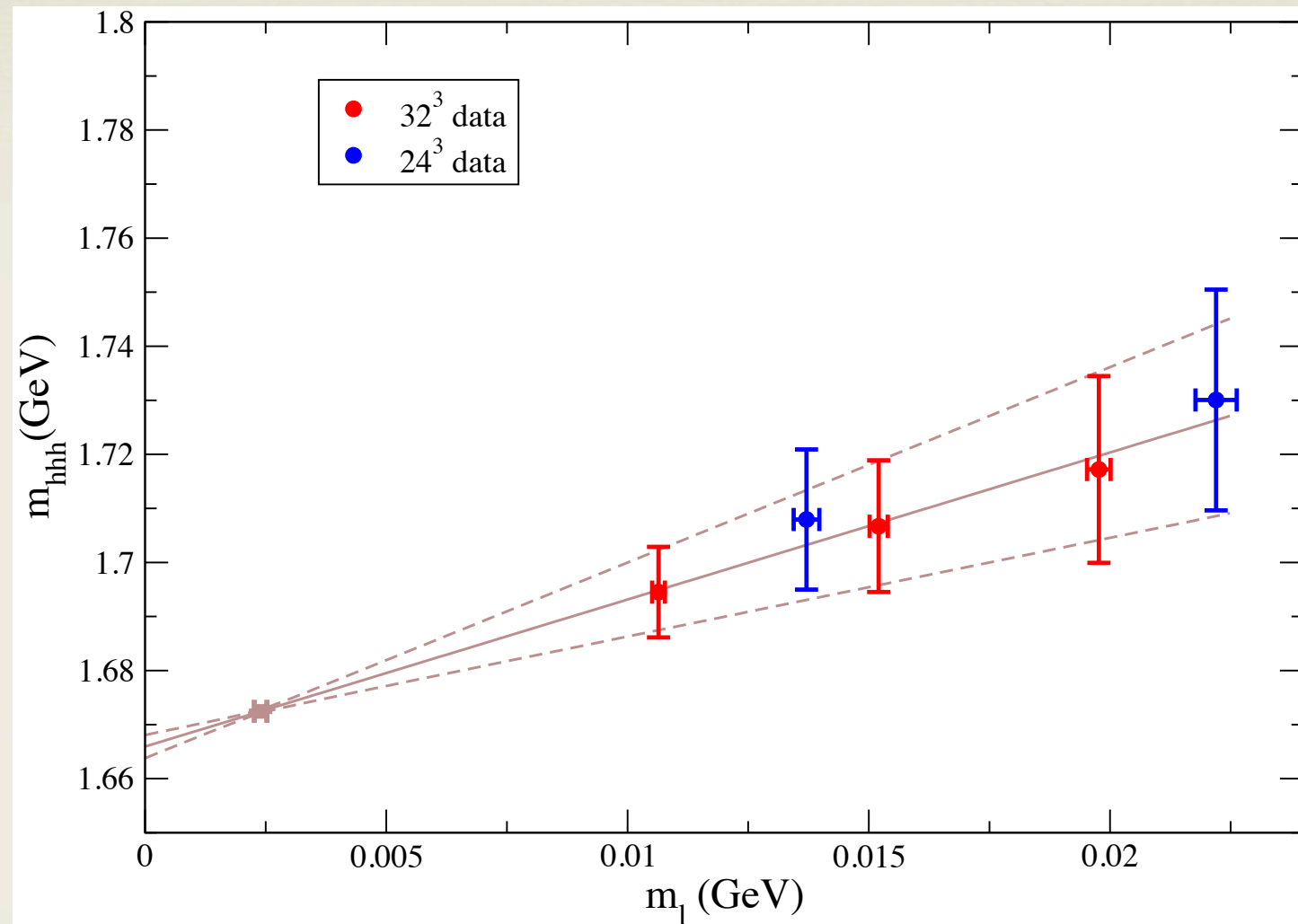


# global fit for non-zero $a$ and $m_{\text{res}}$

- \* NLO power counting:  $a^2$ ,  $m_\pi^2$ : same order, neglect higher order.
- \*  $m_{\text{res}}$  enters as additive renormalization to bare quark mass
- \*  $\tilde{m} = m + m_{\text{res}}$
- \* other  $m_{\text{res}}$  effects are negligible with this power counting
- \* SU(2) ChPT (LEC's depend on  $m_h$ )
- \* Analytic fit to first order in  $m_l$  (constants depend on  $m_h$ )



$m_\Omega$



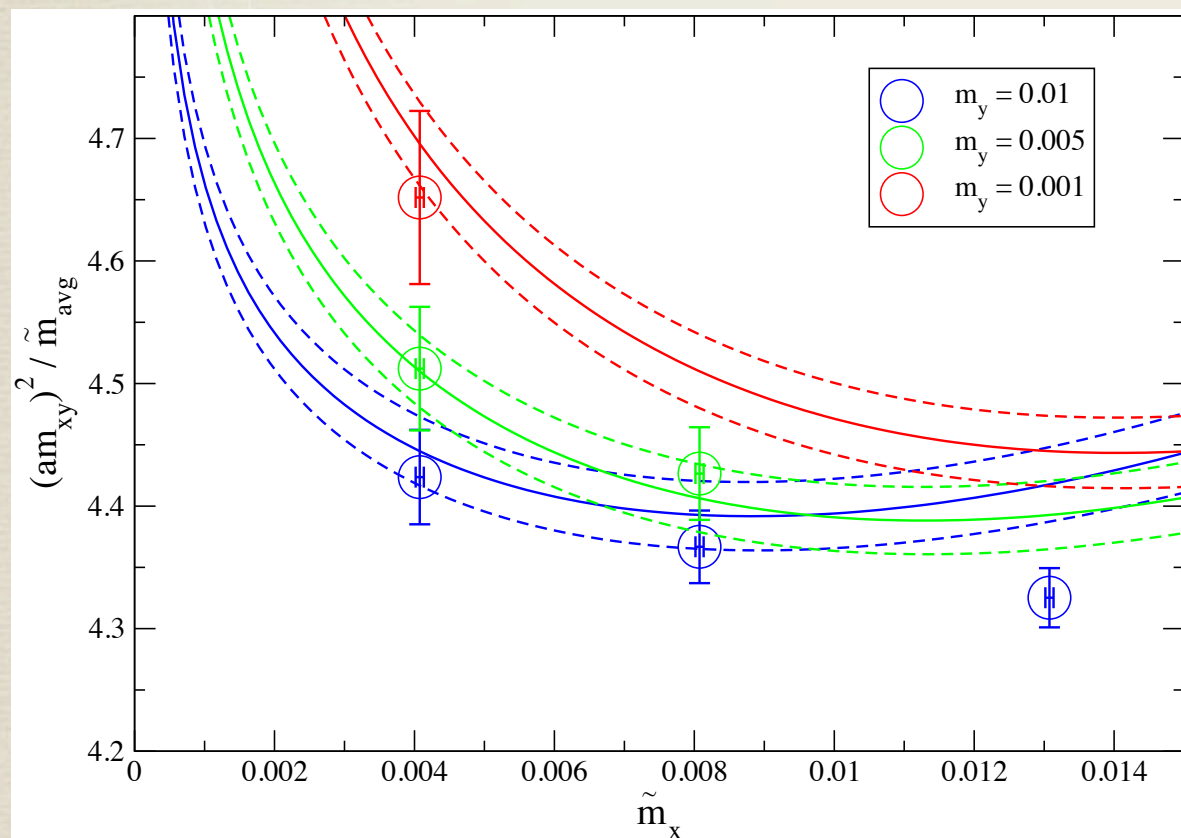
\* analytic ansatz

\* virtually no difference for ChPT and ChPT<sub>FV</sub>

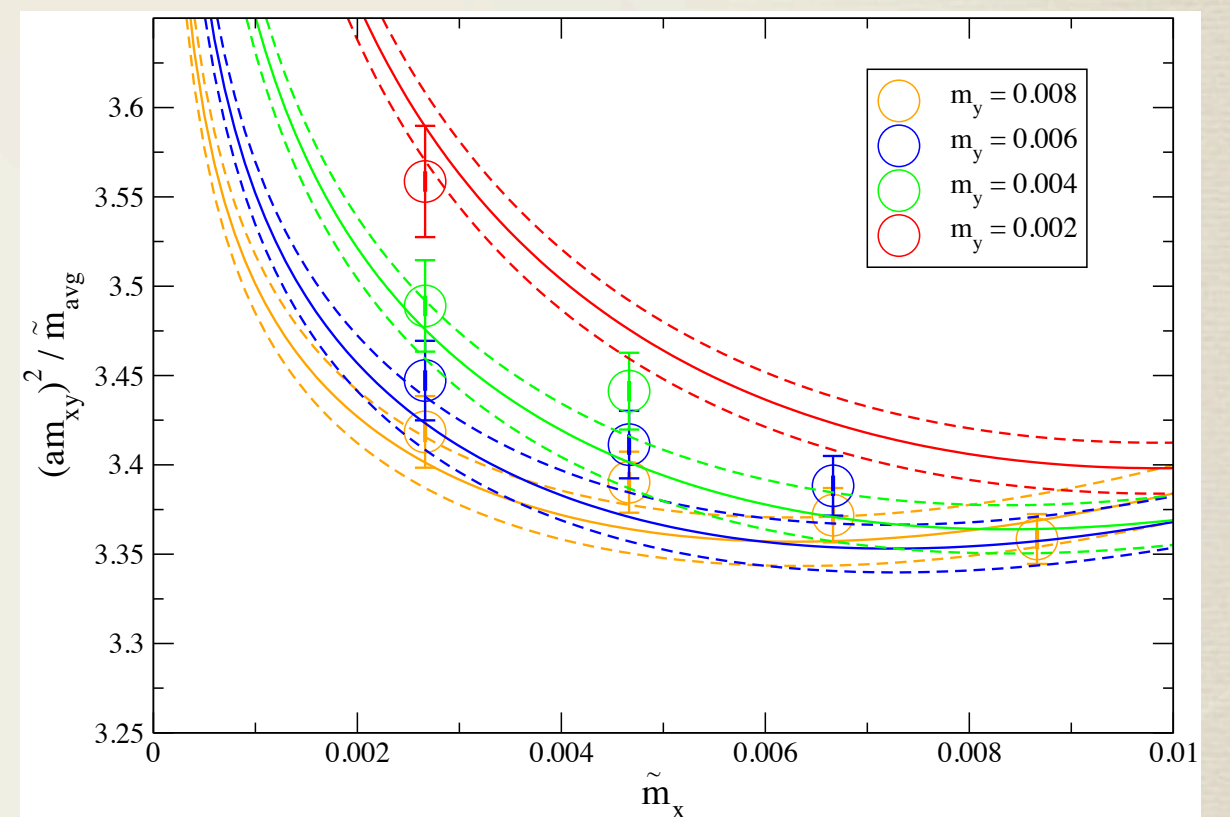


# $m_{xx}^2/m_x$ : partially quenched ChPT<sub>FV</sub> fit

$24^3$ ,  $m_l=0.0005$



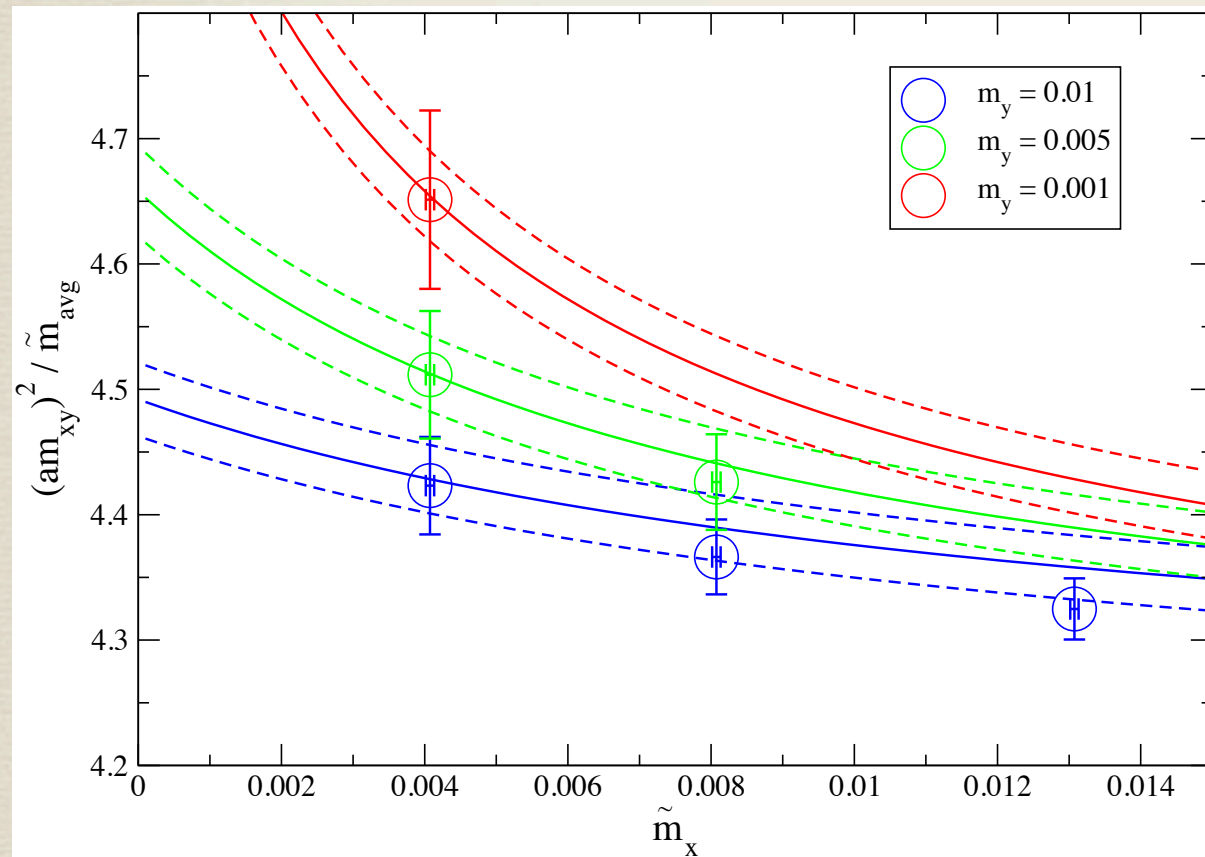
$32^3$ ,  $m_l=0.0004$



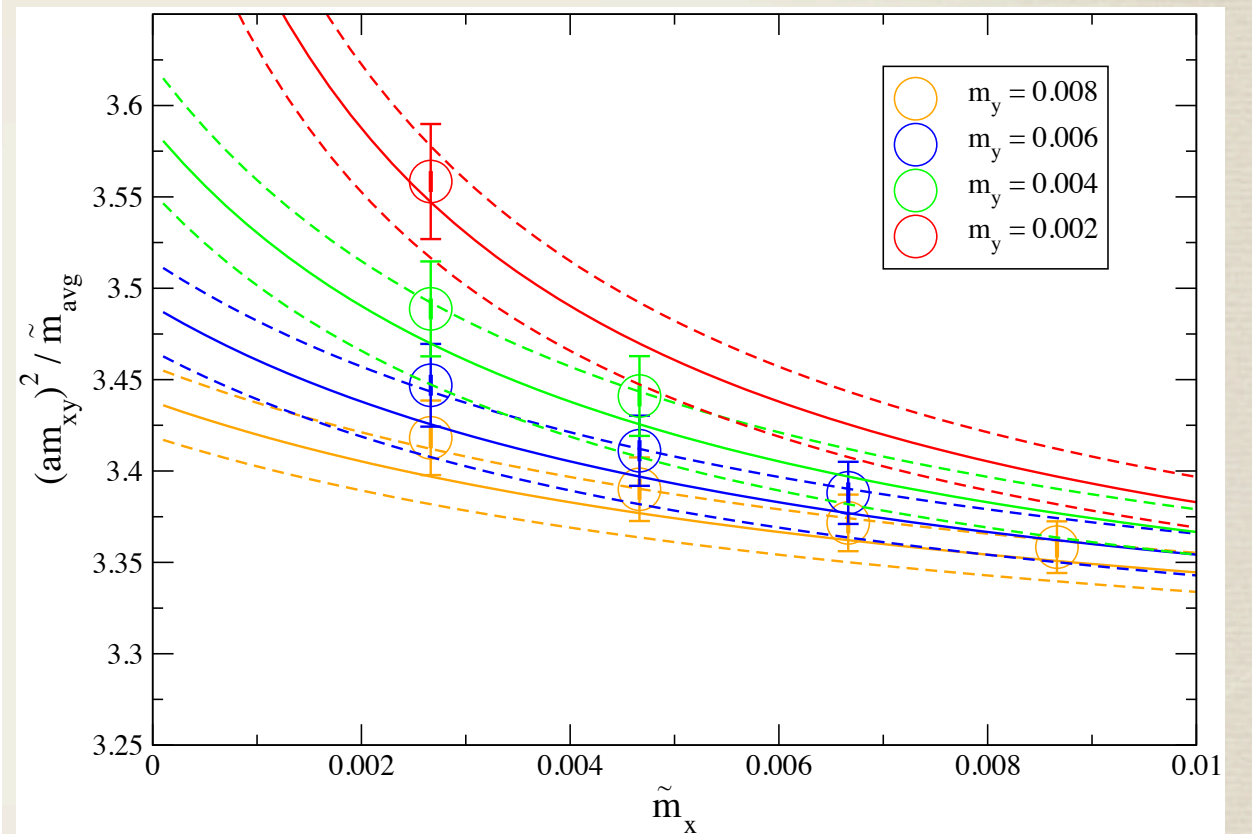


# $m_{xx}^2/m_x$ : analytic fit

$24^3$



$32^3$

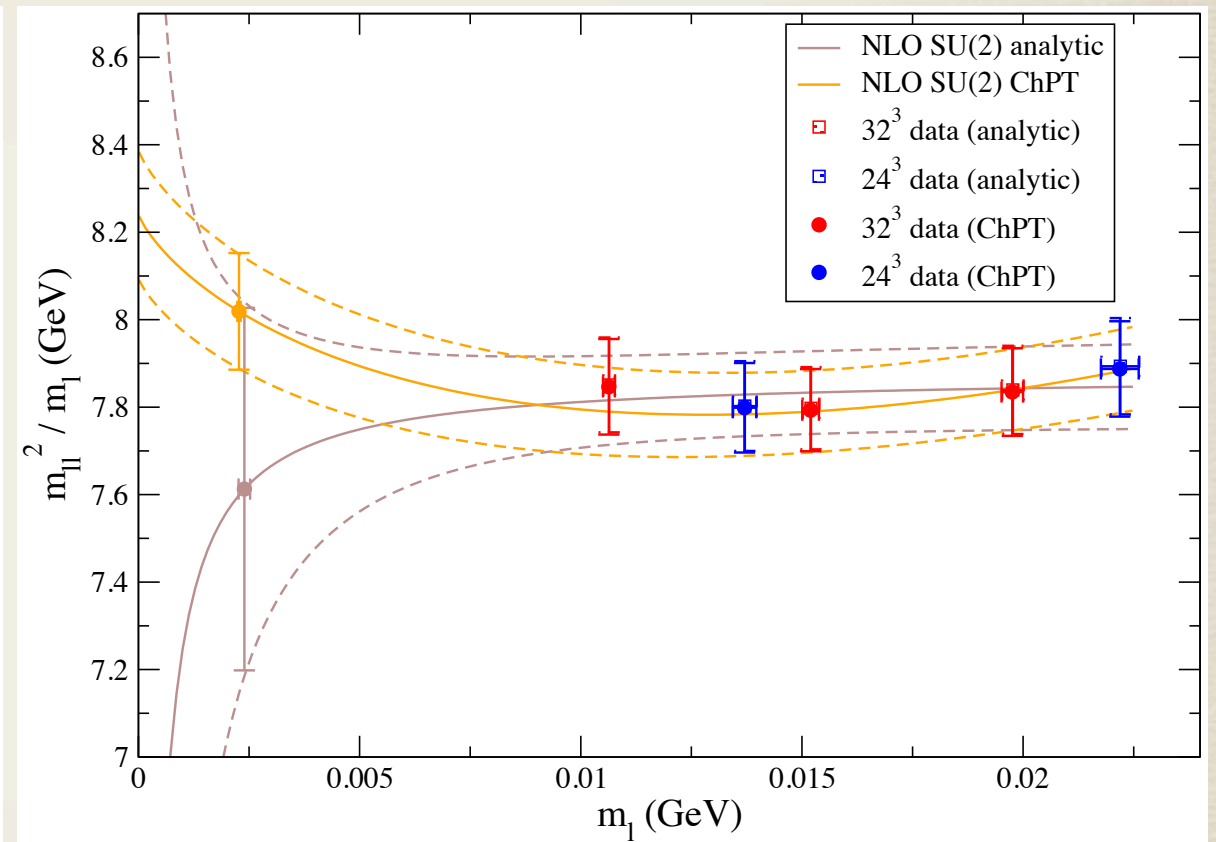
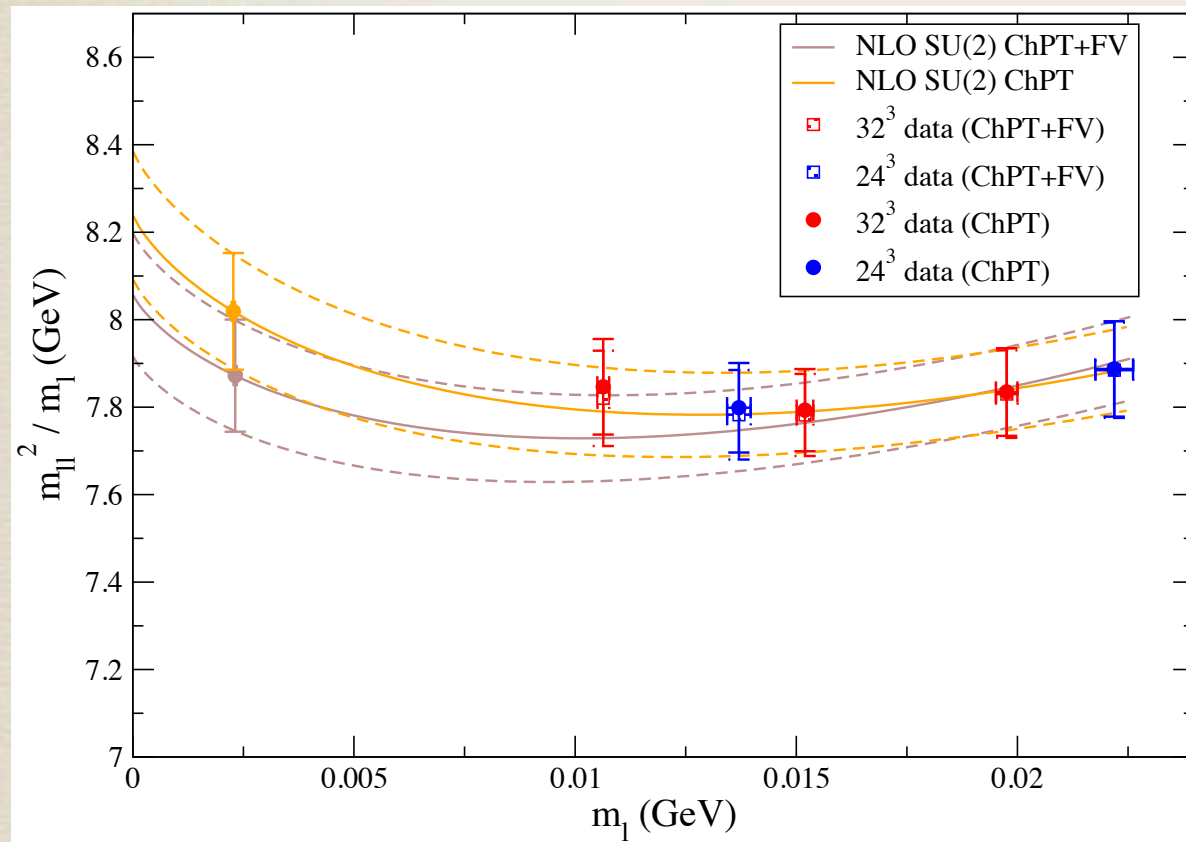




# $m_{\Pi}^2/m_1$ : continuum limit

ChPT( $V \neq \infty$  &  $V \rightarrow \infty$ )

analytic & ChPT( $V \rightarrow \infty$ )



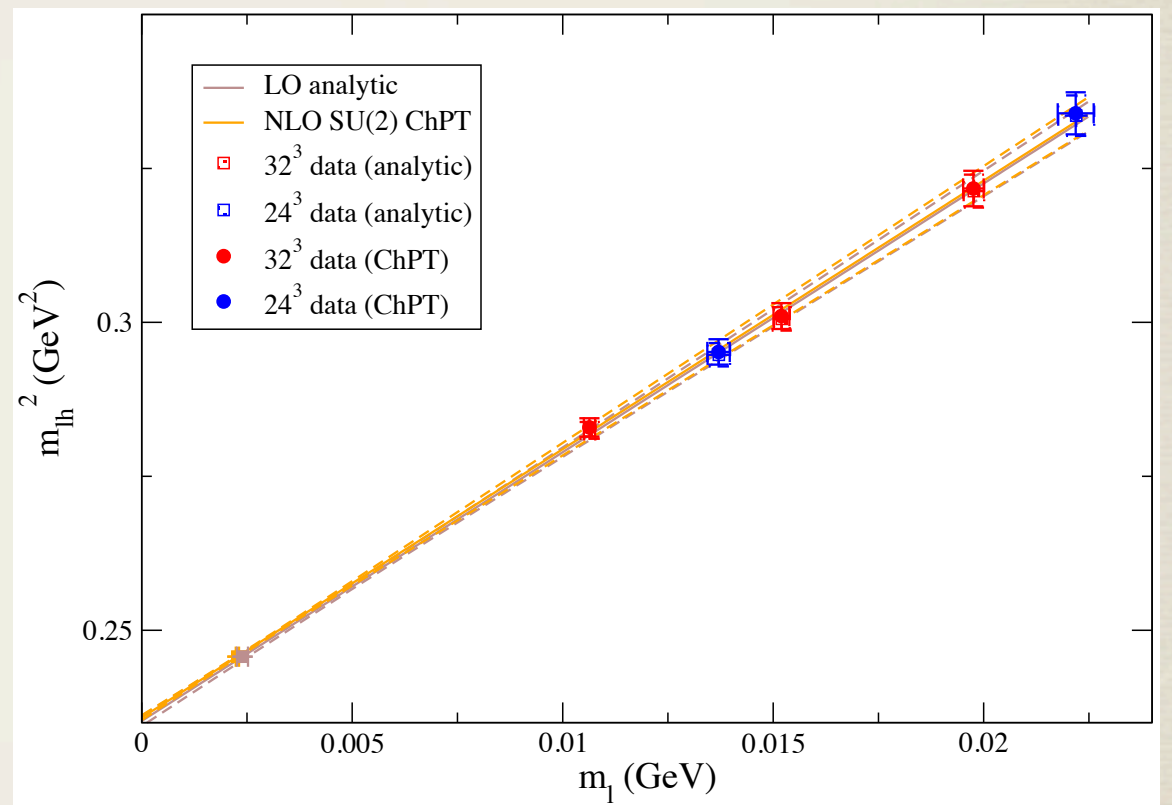
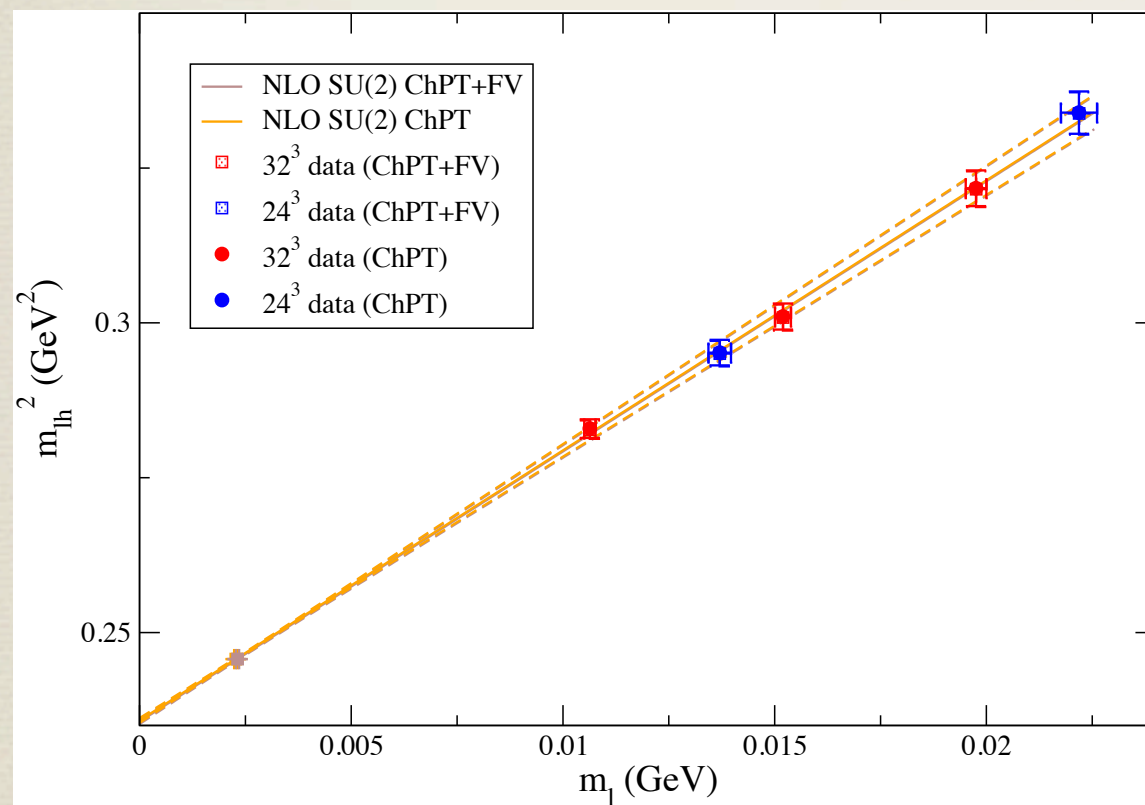
\*  $a^2$  subtracted from the data from the fit



# kaon mass (continuum)

ChPT( $V \neq \infty$  &  $V \rightarrow \infty$ )

analytic & ChPT( $V \rightarrow \infty$ )

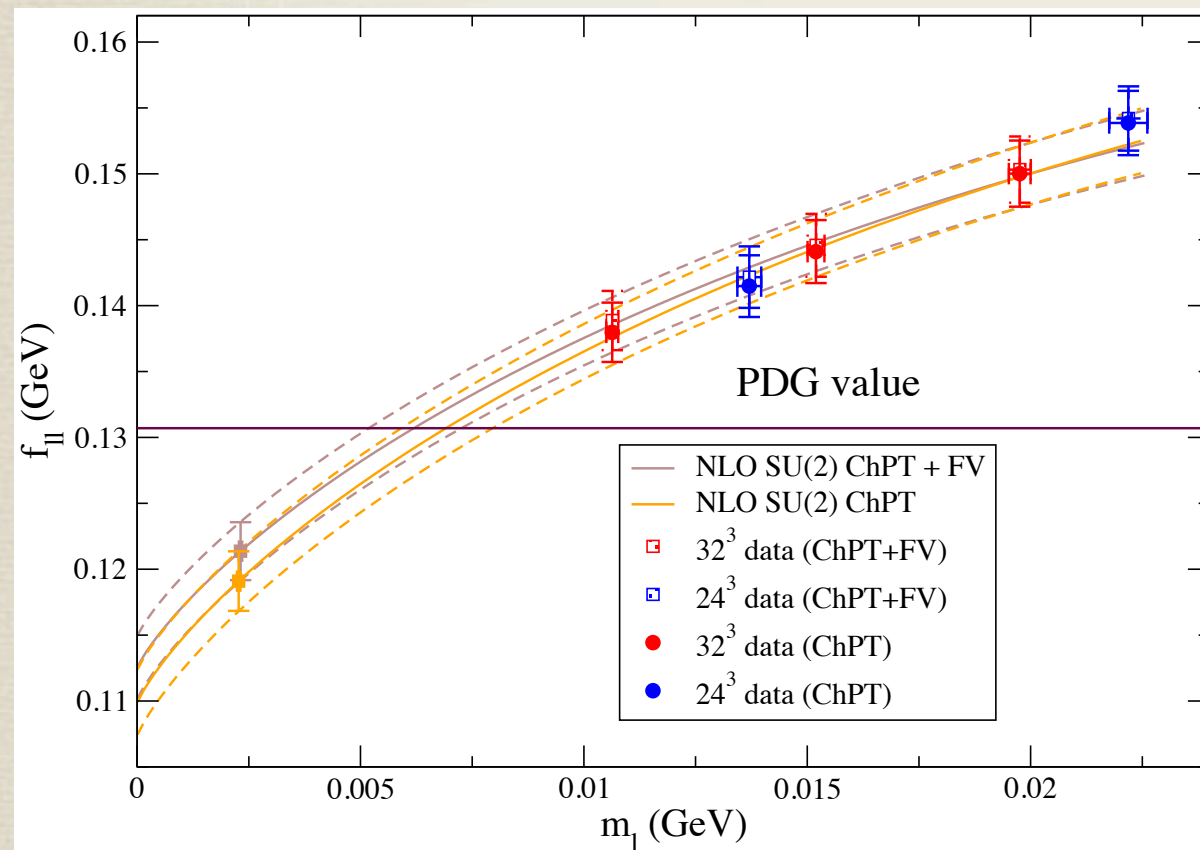


\* local axial current renormalized with  $Z_A$  calculated from the ratio of conserved and local DWF vector current

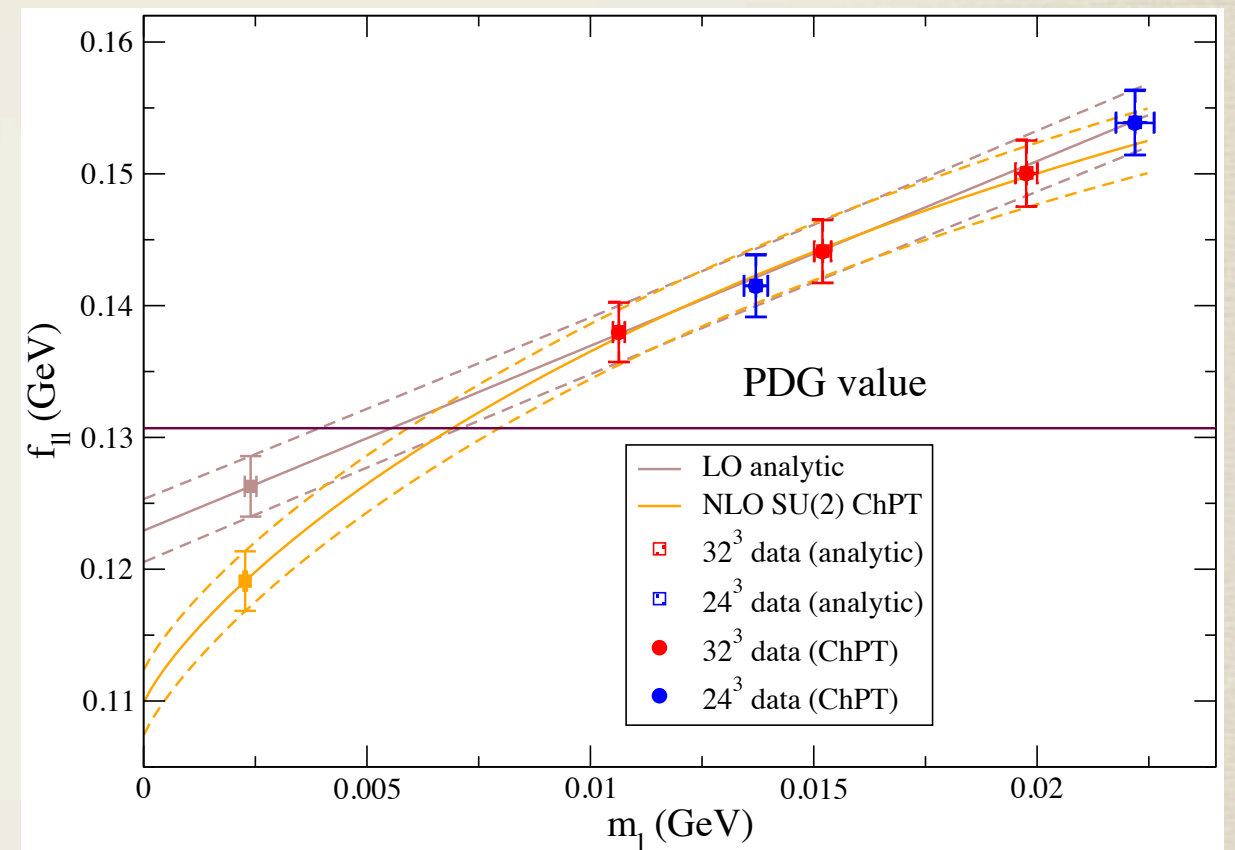


# pion decay constant (continuum)

ChPT( $V \neq \infty$  &  $V \rightarrow \infty$ )



analytic & ChPT( $V \rightarrow \infty$ )



\* local axial current renormalized with  $Z_A$  calculated from the ratio of conserved and local DWF vector current



# central value and systematic error

- \* unitary data : surprisingly linear and consistent with analytic ansatz
- \* but also consistent with ChPT
- \* no strong indication of fit favoring  $\text{ChPT}_{\text{FV}}$  or analytic
- \* take average of them for the central value
- \* systematic error of the chiral extrapolation =  $|\text{ChPT}_{\text{FV}} - \text{analytic}|$
- \* systematic error due to finite volume =  $|\text{ChPT}_{\text{FV}} - \text{ChPT}|$
- \*  $f_{\pi}^{\text{continuum}} = 124(2)(5) \text{ MeV} \leftrightarrow 130.4(4)(2) \text{ MeV [PDG]}$
- \* consistent at 1 sigma level
- \* This procedure is followed for all quantities  $f_K$ ,  $B_K$ , quark masses



# predictions

$$f_{\pi}^{\text{continuum}} = 124(2)(5) \text{ MeV}$$

$$f_K^{\text{continuum}} = 149(2)(4) \text{ MeV}$$

$$(f_K/f_{\pi})^{\text{continuum}} = 1.204(7)(25),$$

$$\tilde{m}_{ud} = 2.35(8)(9) \text{ MeV} \quad \text{and} \quad \tilde{m}_s = 63.7(9)(1) \text{ MeV}$$

in 32 cubed scheme. needs to be matched  
to more convenient schemes with non-  
perturbative renormalization



# non-perturbative renormalization

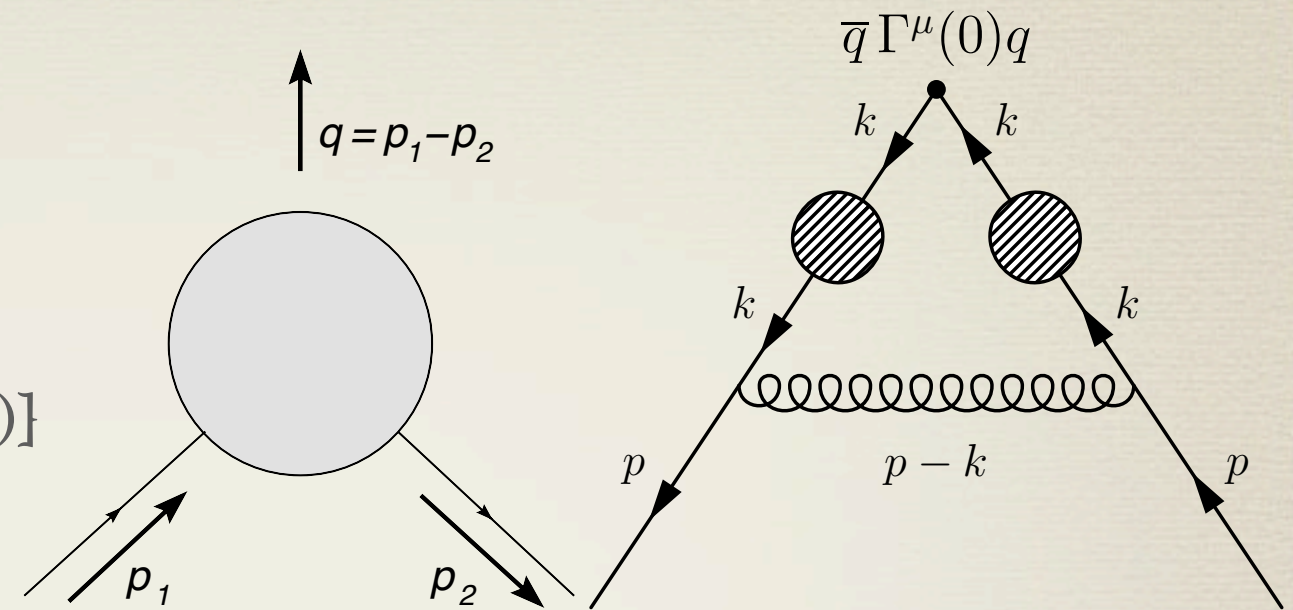
- \* previous RBC/UKQCD result [PRD78(2008)114509]

$$m_s^{\overline{MS}}(2GeV) = 107.33(4.4)_{\text{stat}}(4.9)_{\text{syst}}(9.7)_{\text{ren}} \text{ MeV}$$

- \* **9%** error from RI/MOM renormalization [PRD78(2008)054510]
  - \* 6% from truncation error of perturbation (NNNLO)
  - \* 7% from  $m_s \neq 0$  (SSB contamination)
    - \* a promising solution also provided in the same paper
      - \* momentum kinematics: exceptional  $\rightarrow$  non-exceptional
- \* RI/SMOM schemes are constructed utilizing non-exceptional momenta:
  - \* [C. Stuttmann, Y. A. N. Christ, T. Izubuchi, C. Sachrajda, A. Soni, PRD80(2009)014501]
  - \* 1 loop matching from SMOM schemes to  $\overline{MS}$  provided



# SSB contamination



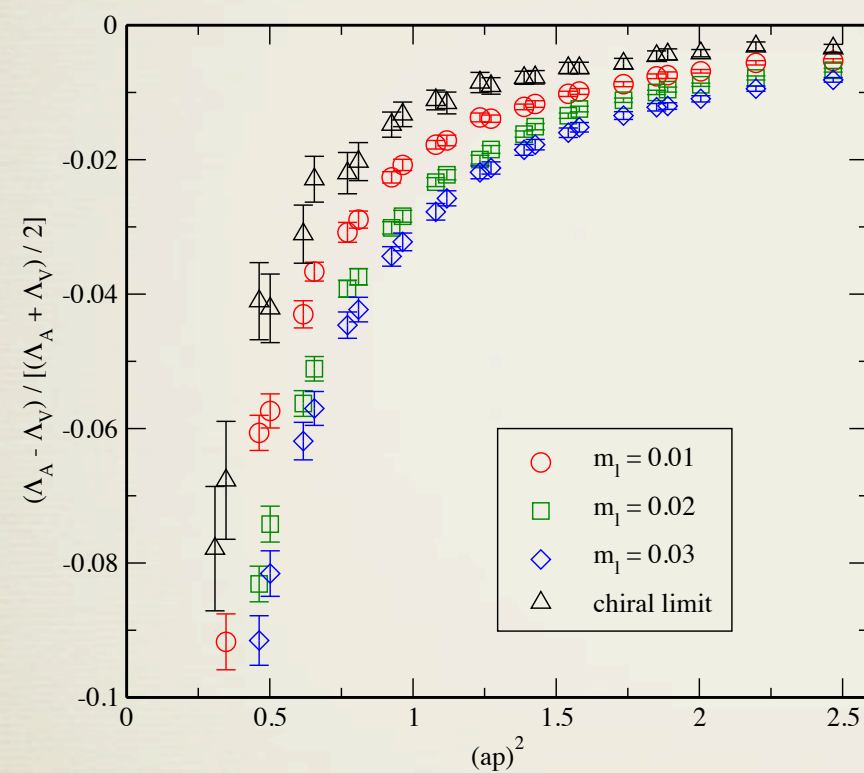
- \*  $q_\mu = 0$  [exceptional ( $\exists$  partial sum is zero)]
- \*  $1/p^2$  from one gluon exchange
- \* low momentum flow in the upper part triggers SSB depending on the  $\Gamma$  structure  
 $\rightarrow 1/p^2$  contamination (cannot be corrected by PT)
- \* Backed up from power counting theorem by Weinberg, combined with the group theoretical argument (RBC/UKQCD, PRD 2008)
- \* avoiding  $q_\mu = 0$ , and going for non-exceptional momenta remove  $1/p^2$  contamination



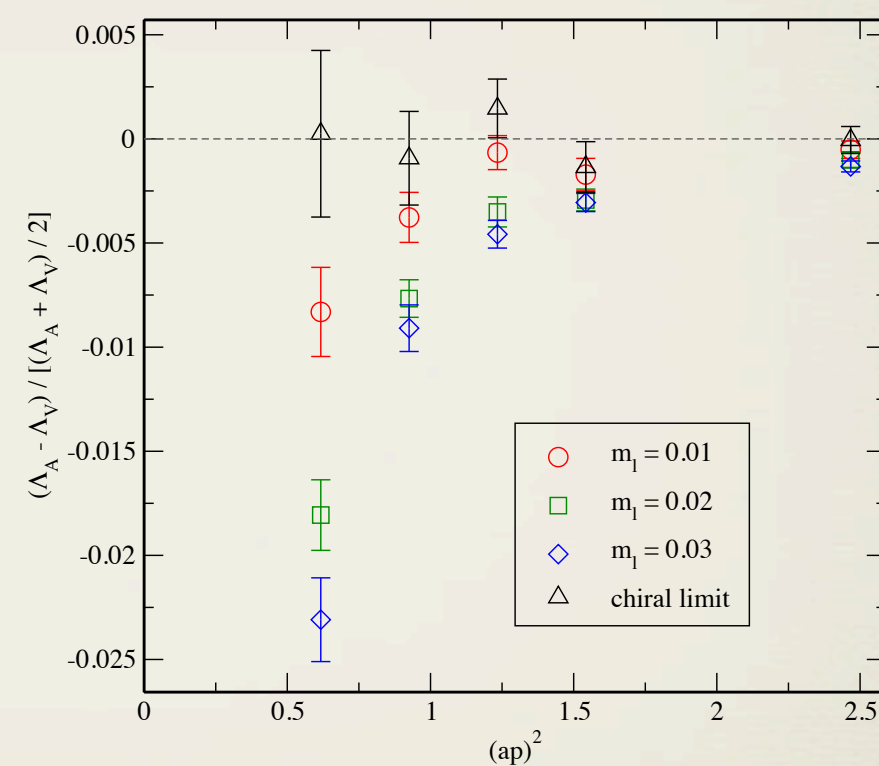
# A test of non-exceptional mom

$\Lambda_A - \Lambda_V$ : RBC/UKQCD [PRD 2008]

exceptional



non-exceptional

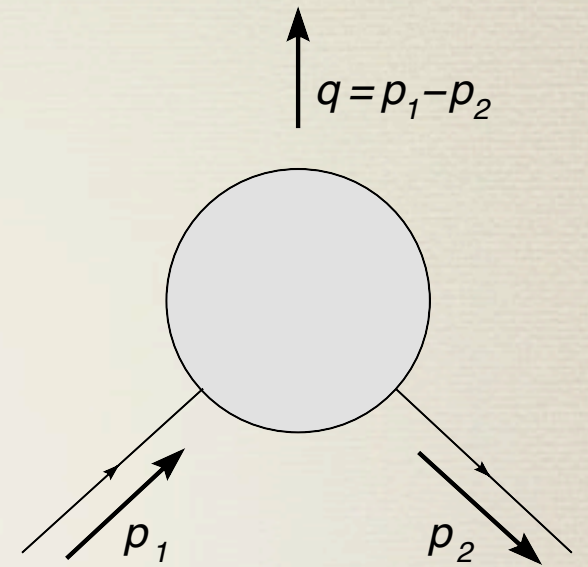


\* The success created a very good motivation to invest in non-exceptional momenta



# SMOM scheme

- \* C. Sturm et al. PRD80 (2009) 014501.
- \* utilize a **non-exceptional momenta** for bilinears
- \* Symmetric ( $q^2 = p_1^2 = p_2^2$ ) MOM scheme
- \* SSB contamination is expected to be reduced
- \* SMOM:  $\Lambda_A = \Lambda_V$  exact for chiral fermions
- \*  $\Lambda_A = \Lambda_V + c/p^2 \dots$  for RI/MOM
- \* projection operator for V:  $q_\mu \not{q} / q^2$  gives  $Z_q$  of RI'
- \* 2nd scheme: SMOM $\gamma_\mu$   $\gamma_\mu$  new  $Z_q$





# MOM $\leftrightarrow$ SMOM comparison: $\overline{MS}$ matching

\* original RI/MOM mass conversion factor

[Franco & Lubicz, Chetyrkin & Retey, Gracey]

$\mu = 2 \text{ GeV}$

$$C_m(RI/MOM \rightarrow \overline{MS}) = 1 + \frac{\alpha_s}{4\pi} C_F \frac{4}{4} + \dots = 1 - \textcolor{red}{0.123} - \textcolor{blue}{0.070} - 0.048 + \dots$$

1-loop, 2-loop, 3-loop

$$C_m(RI'/MOM \rightarrow \overline{MS}) = 1 + \frac{\alpha_s}{4\pi} C_F \frac{4}{4} + \dots = 1 - \textcolor{red}{0.123} - \textcolor{blue}{0.065} - 0.044 + \dots$$

5-6% correction 2  $\rightarrow$  3 loop!

\* SMOM: 1 loop [C. Sturm et al. PRD80 (2009) 014501]

\* 2 loop [L. Almeida & C. Sturm PRD82 (2010) 054017]

$$C_m(SMOM \rightarrow \overline{MS}) = 1 - \frac{\alpha_s}{4\pi} C_F \frac{0.484}{4} + \dots = 1 - \textcolor{red}{0.015} - \textcolor{blue}{0.006} + \dots$$

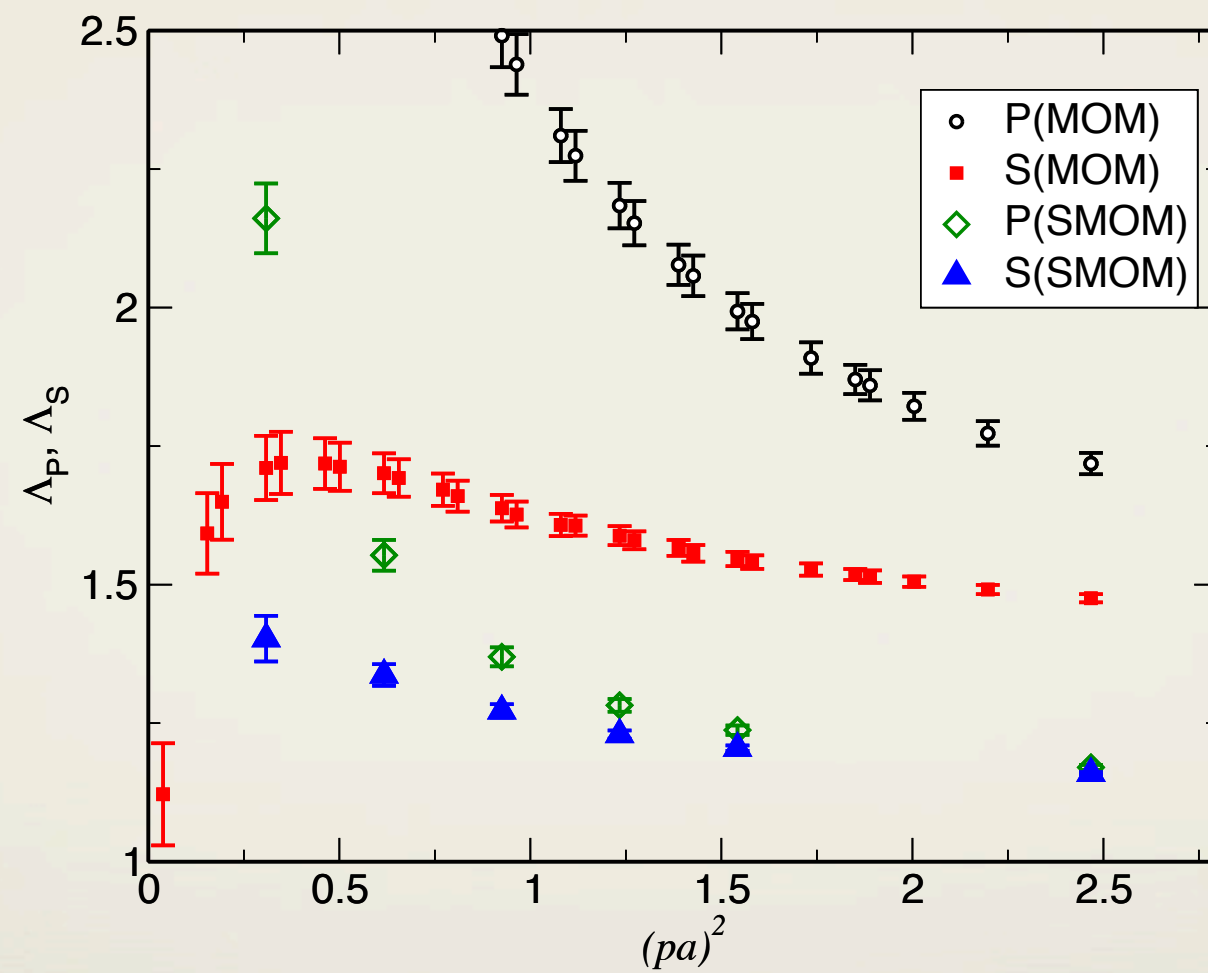
$$C_m(SMOM\gamma_\mu \rightarrow \overline{MS}) = 1 - \frac{\alpha_s}{4\pi} C_F \frac{1.484}{4} + \dots = 1 - \textcolor{red}{0.045} - \textcolor{blue}{0.020} + \dots$$

\* (size of the last term is taken as the systematic error of PT)



# results with 2+1f DWF

$\Lambda_P, \Lambda_S$  comparison of RI/MOM and SMOM

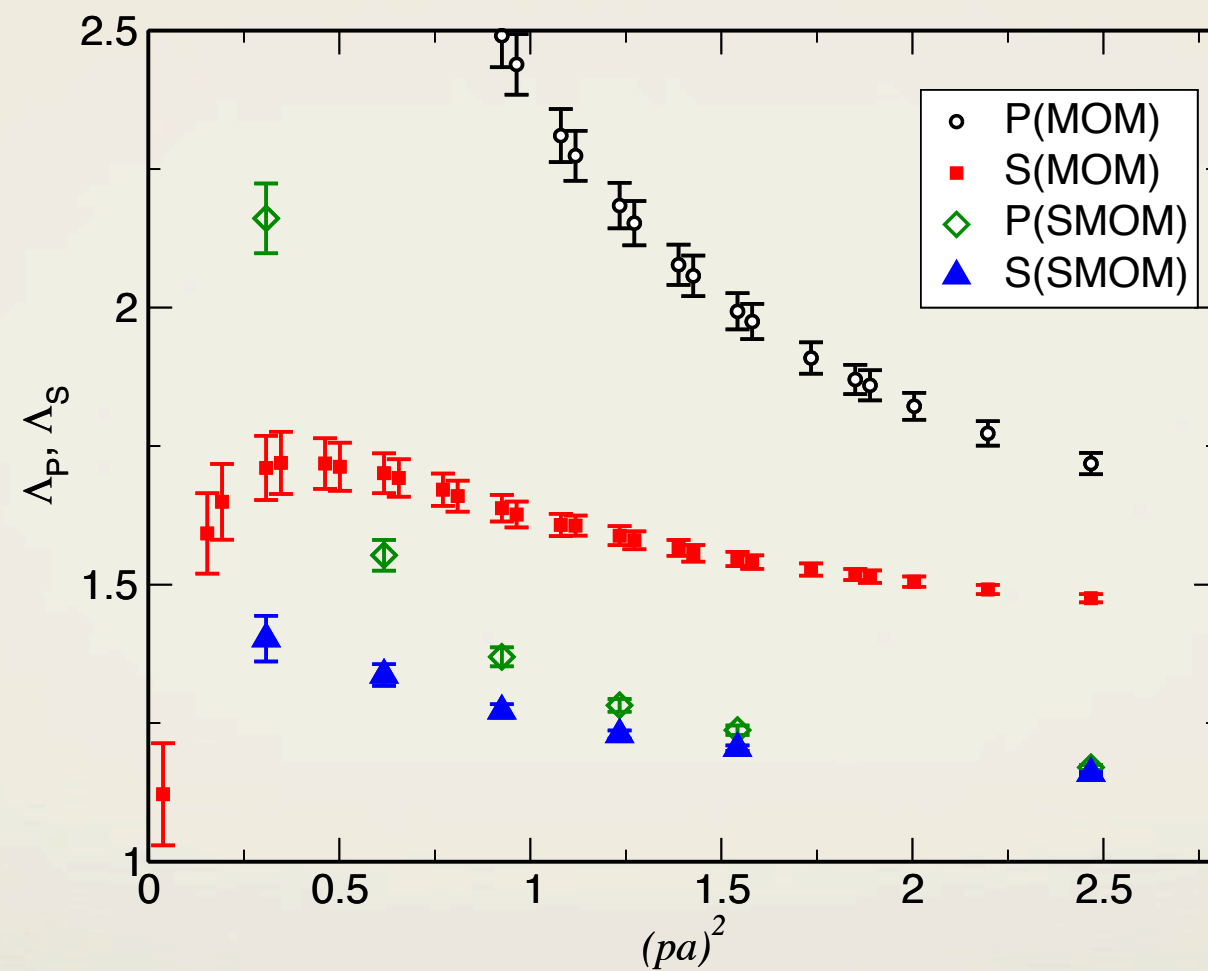




# results with 2+1f DWF

## $\Lambda_P, \Lambda_S$ comparison of RI/MOM and SMOM

\* DWF, RBC/UKQCD (YA:Lattice 2008)

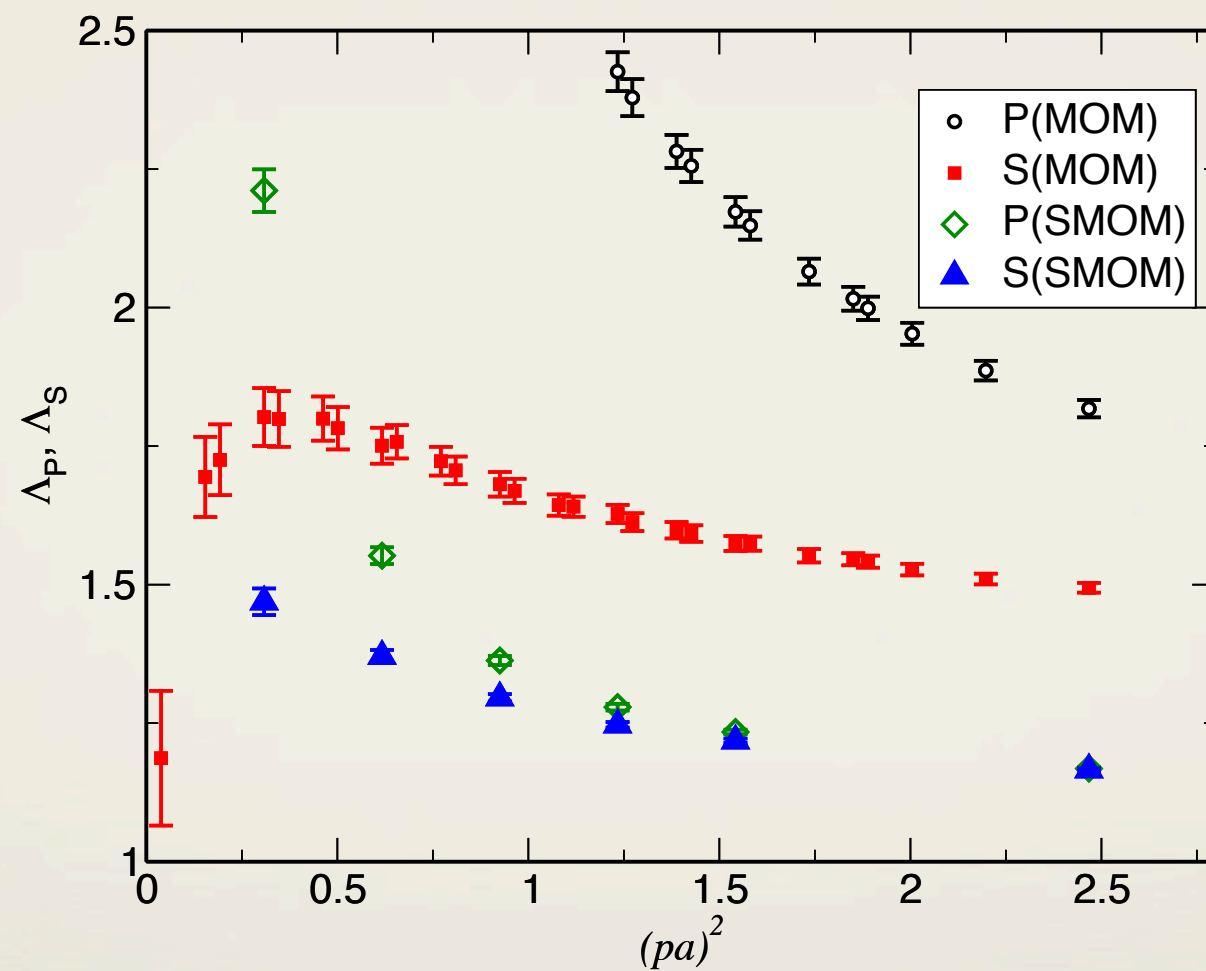




# results with 2+1f DWF

## $\Lambda_P, \Lambda_S$ comparison of RI/MOM and SMOM

\* DWF, RBC/UKQCD (YA:Lattice 2008)

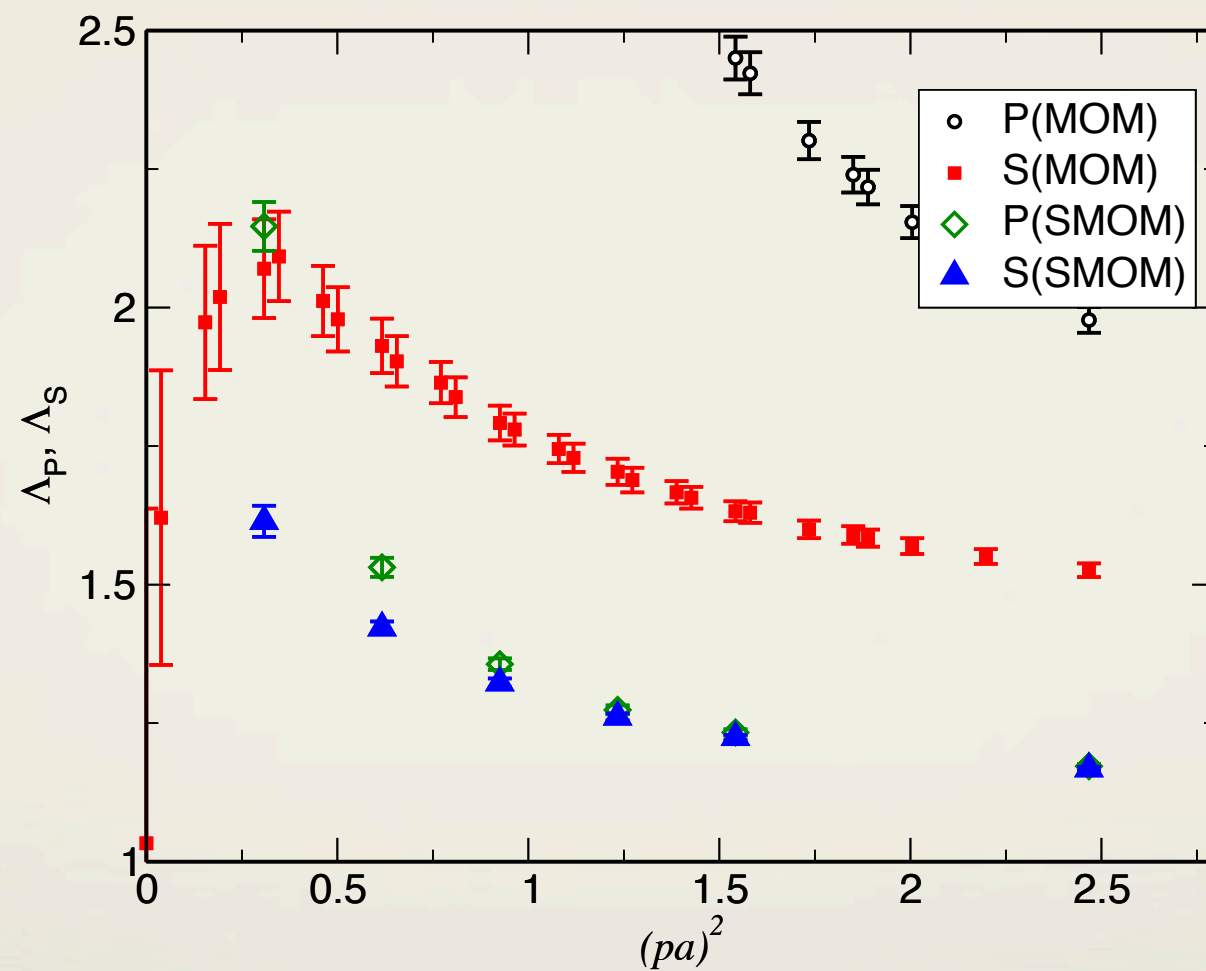




# results with 2+1f DWF

## $\Lambda_P, \Lambda_S$ comparison of RI/MOM and SMOM

\* DWF, RBC/UKQCD (YA:Lattice 2008)

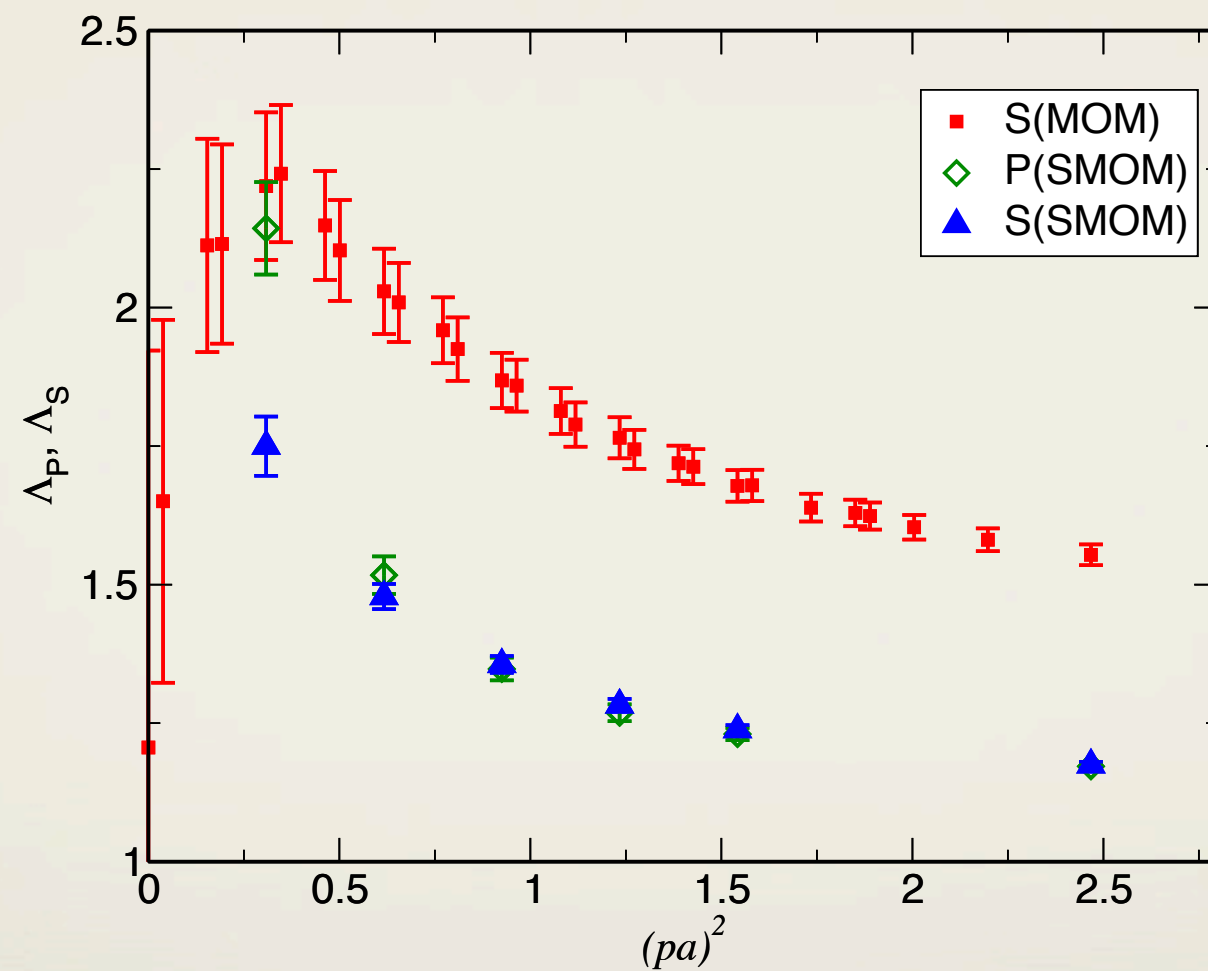




# results with 2+1f DWF

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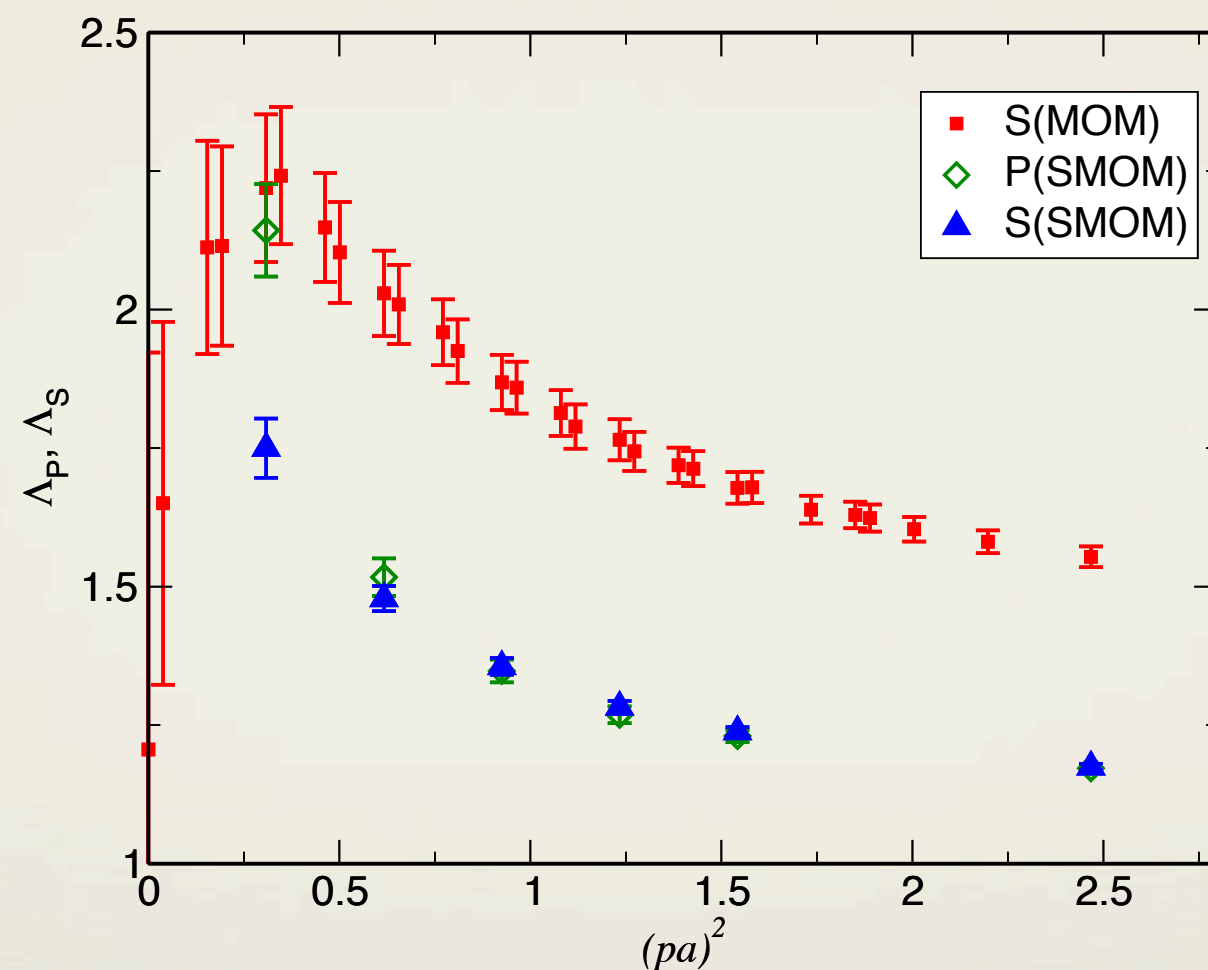




# results with 2+1f DWF

## $\Lambda_P, \Lambda_S$ comparison of RI/MOM and SMOM

\* DWF, RBC/UKQCD (YA:Lattice 2008)



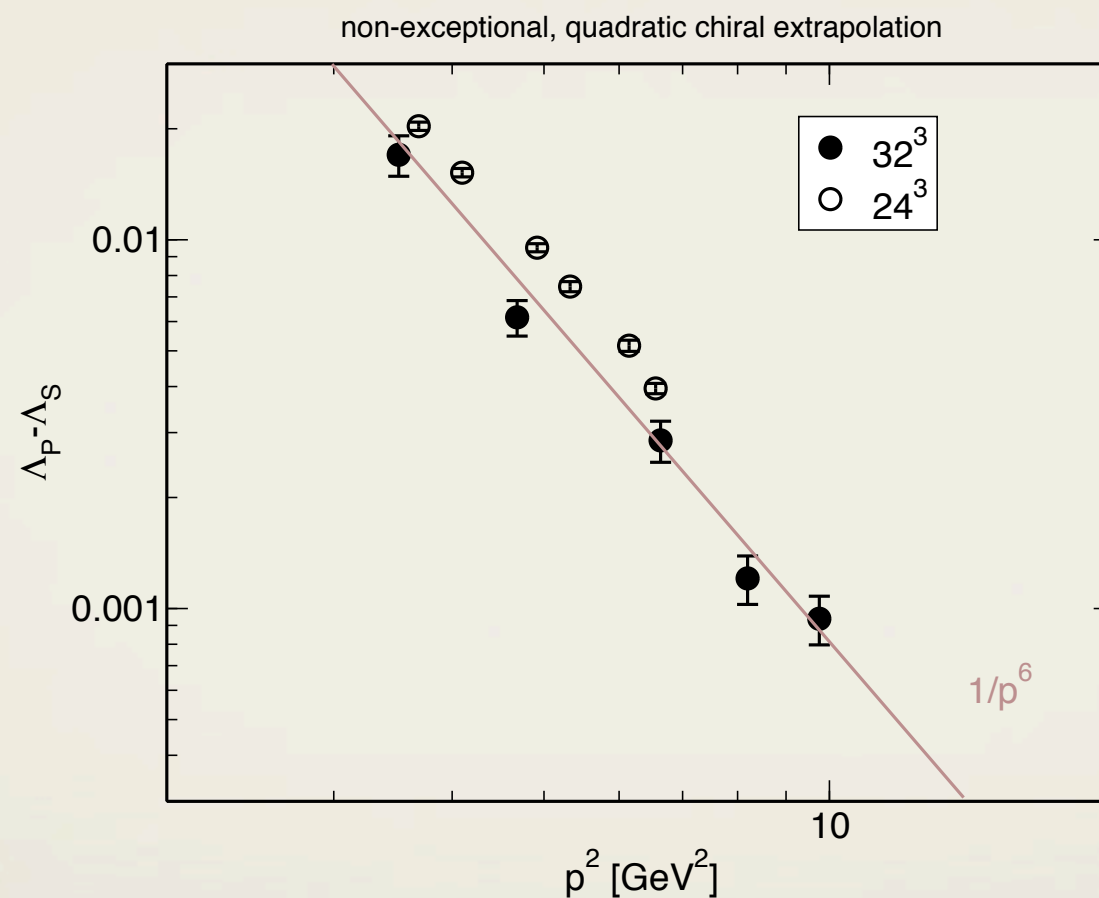
\* SMOM: smaller mass dependence, good chiral symm.



# SMOM $\Lambda_P - \Lambda_S$

on two lattice spacings, with powerful momentum source

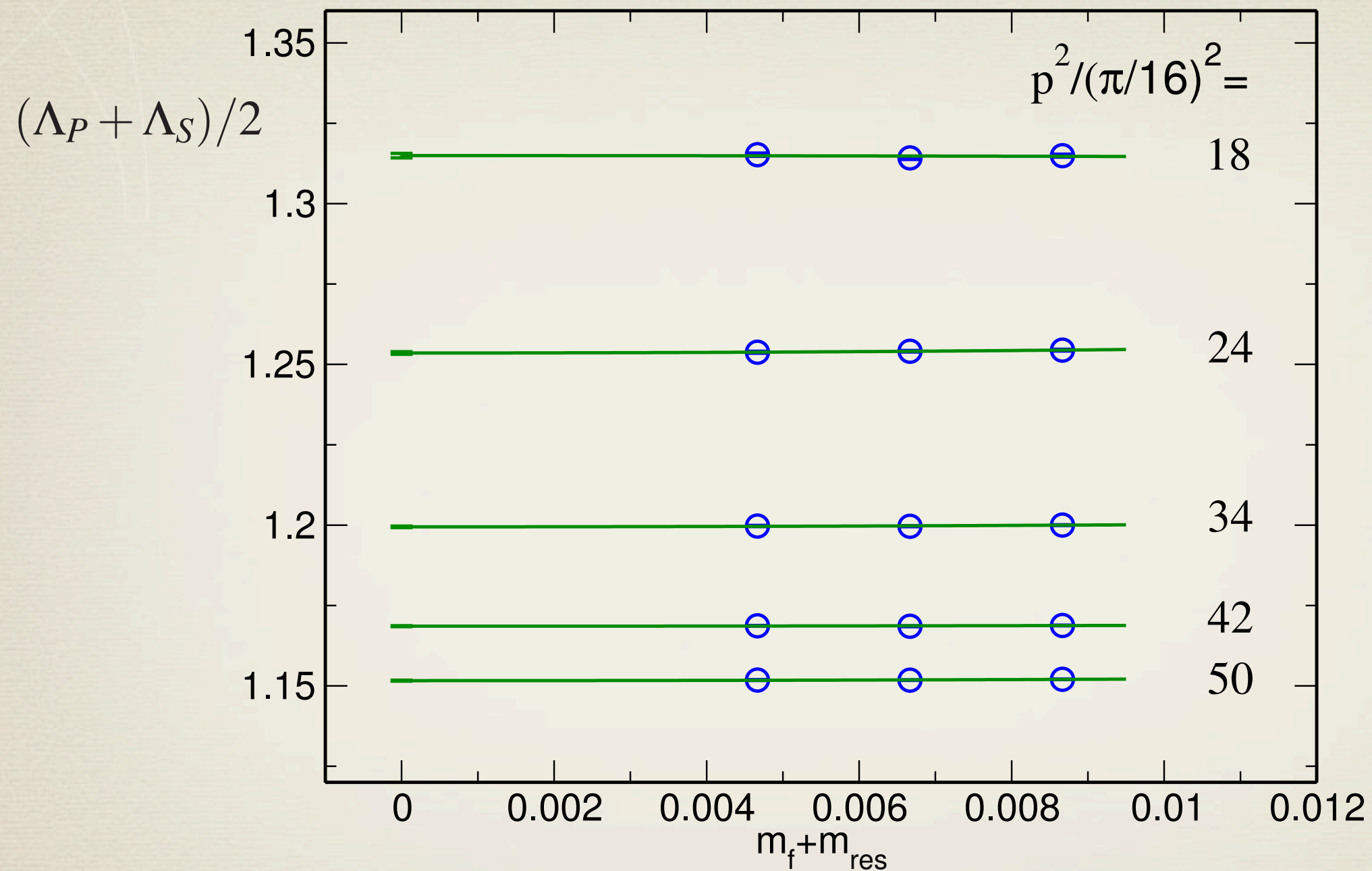
\*  $a^{-1} = 1.7\text{GeV}$  ( $24^3$  coarse) &  $2.3\text{GeV}$  ( $32^3$  fine), DWF



\*  $1/p^6$  expected from the Weinberg's theorem



# SMOM kinematics: small mass dependence

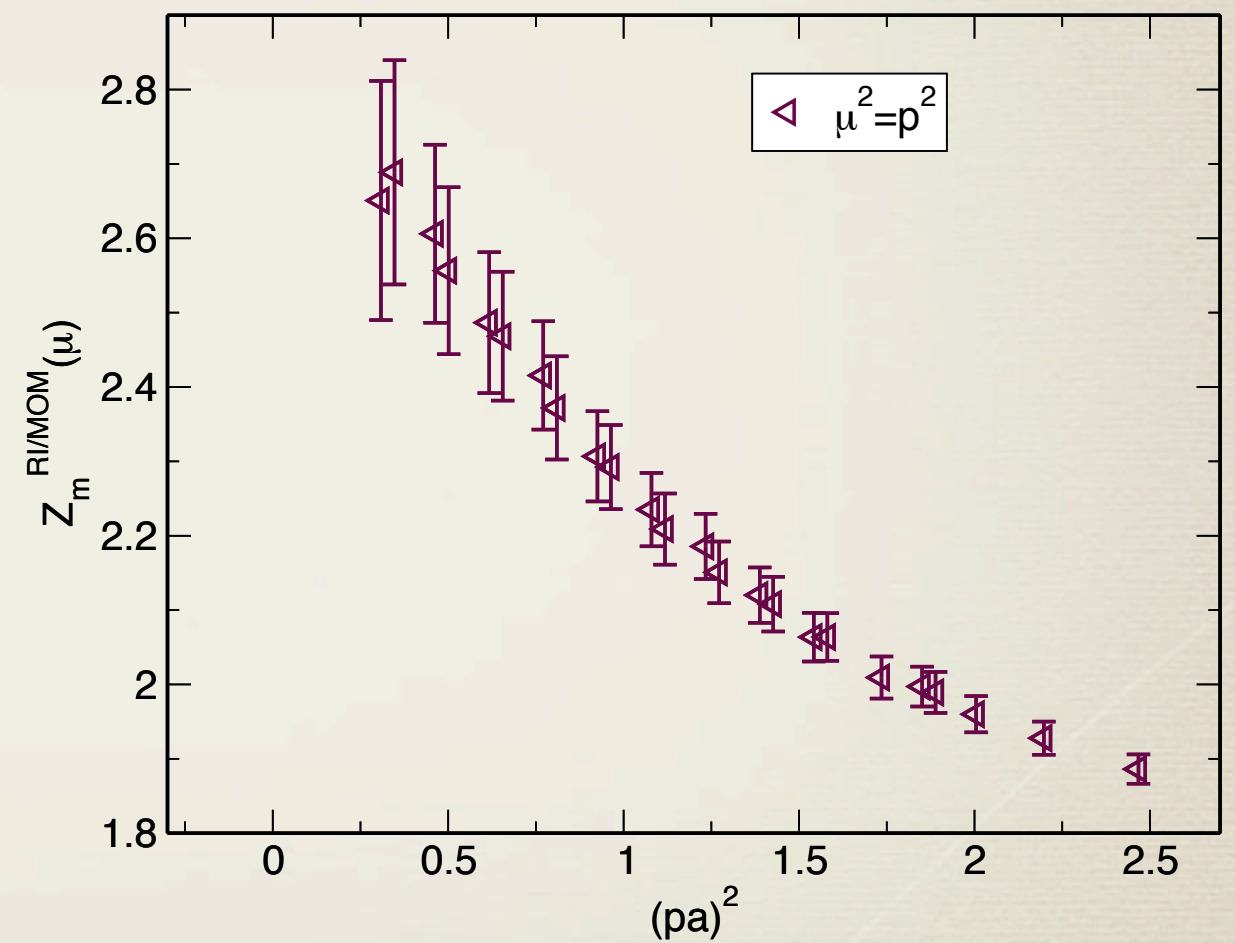


\*  $m_s \neq 0$  effect is estimated as 0.1-0.2 % level



# RI/MOM procedure

RBC/UKQCD  
 $n_f=2+1$  DWF,  $1/a \approx 1.7$  GeV

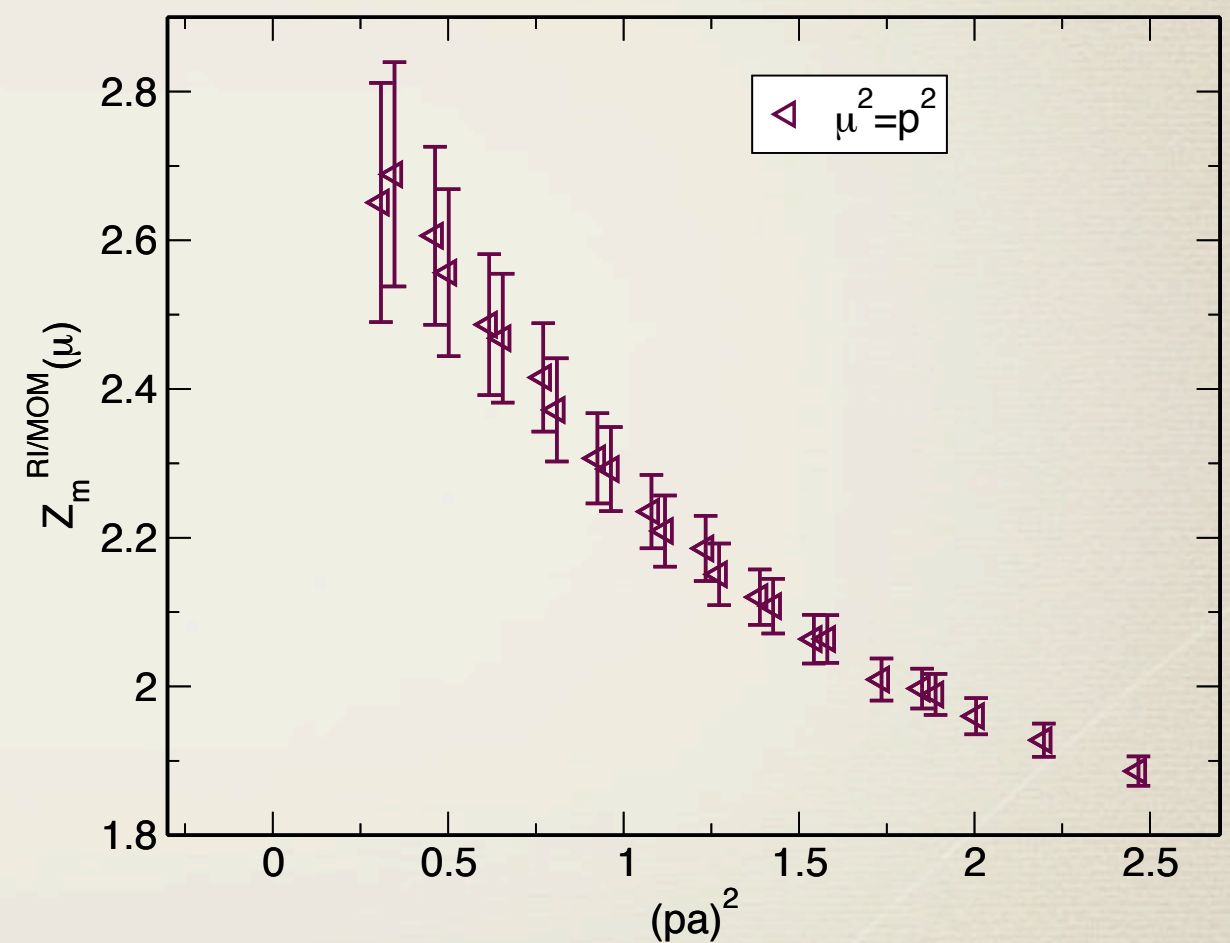




# RI/MOM procedure

\*  $Z_m^{\text{RI}}(\mu)$  non-perturbative  $\mu$  dependence

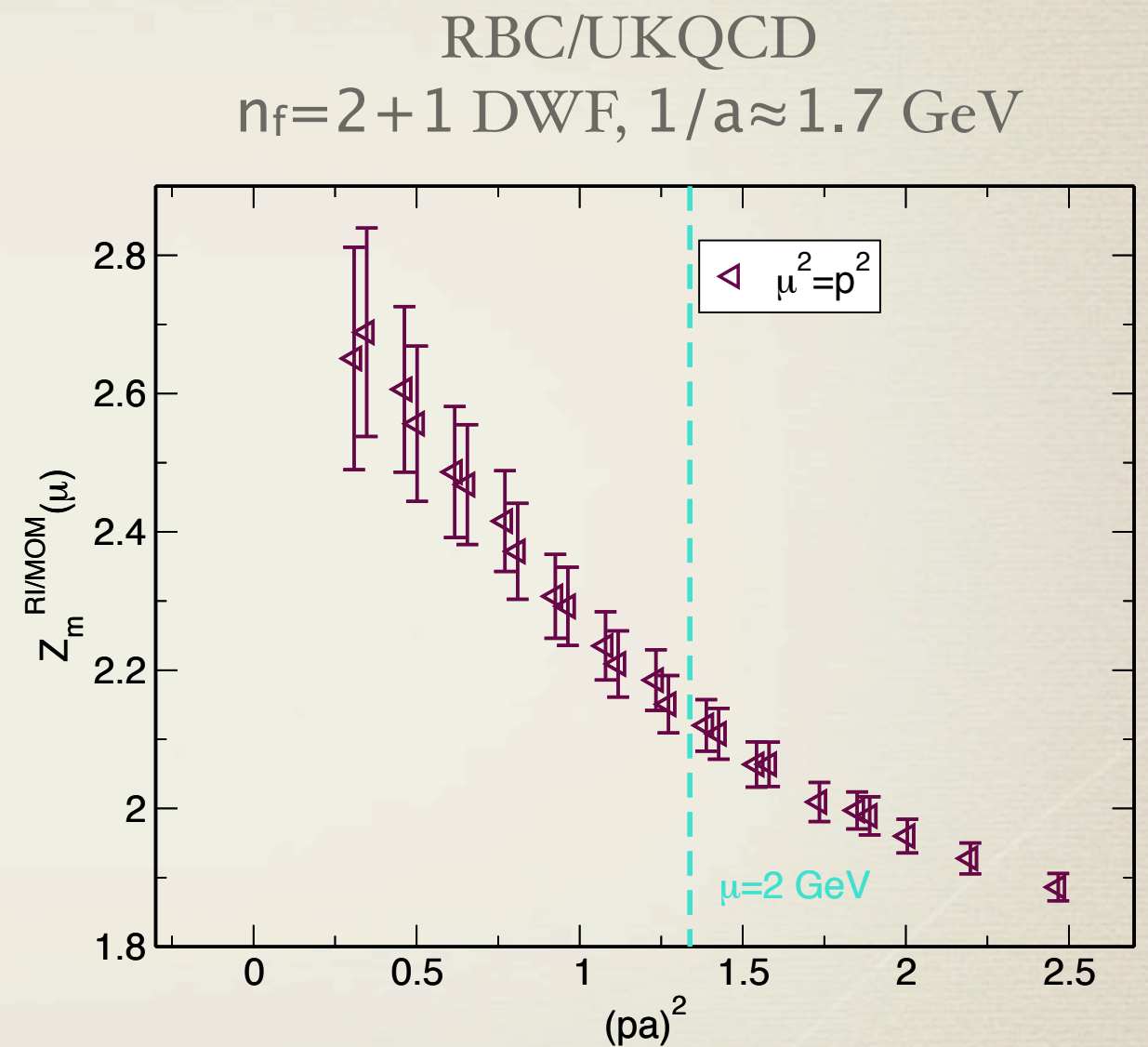
RBC/UKQCD  
 $n_f=2+1$  DWF,  $1/a \approx 1.7$  GeV





# RI/MOM procedure

- \*  $Z_m^{\text{RI}}(\mu)$  non-perturbative  $\mu$  dependence
- \*  $Z_m^{\text{RI}}(2 \text{ GeV})$  from  $p=2 \text{ GeV}$  intercept



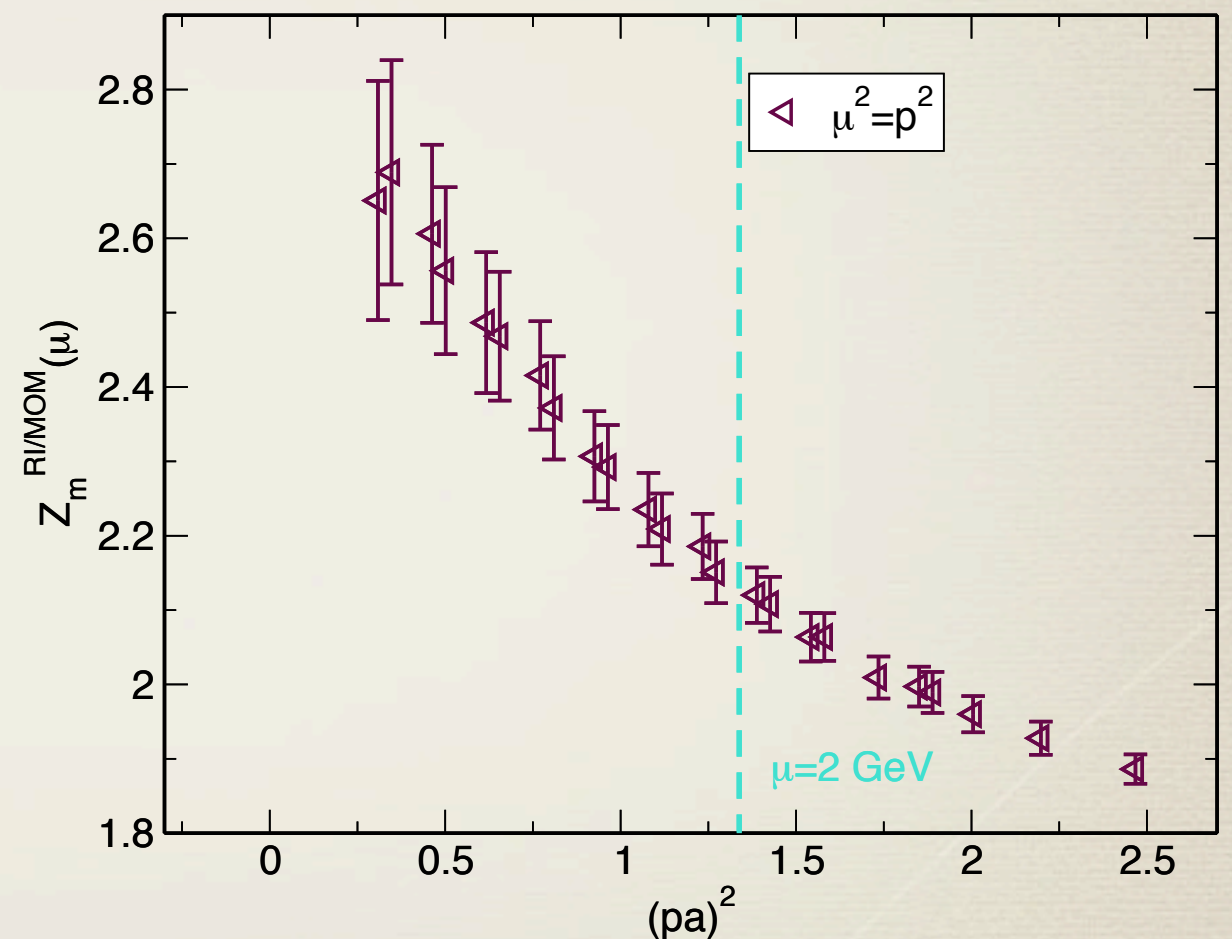


# RI/MOM procedure

- \*  $Z_m^{\text{RI}}(\mu)$  non-perturbative  $\mu$  dependence
- \*  $Z_m^{\text{RI}}(2 \text{ GeV})$  from  $p=2 \text{ GeV}$  intercept

$$Z_m^{\overline{MS}}(\mu) = \left( \frac{Z_m^{\overline{MS}}(\mu)}{Z_m^{\text{RI}}(\mu)} \right)_{PT} Z_m^{\text{RI,NPR}}(\mu)$$

RBC/UKQCD  
 $n_f=2+1$  DWF,  $1/a \approx 1.7 \text{ GeV}$



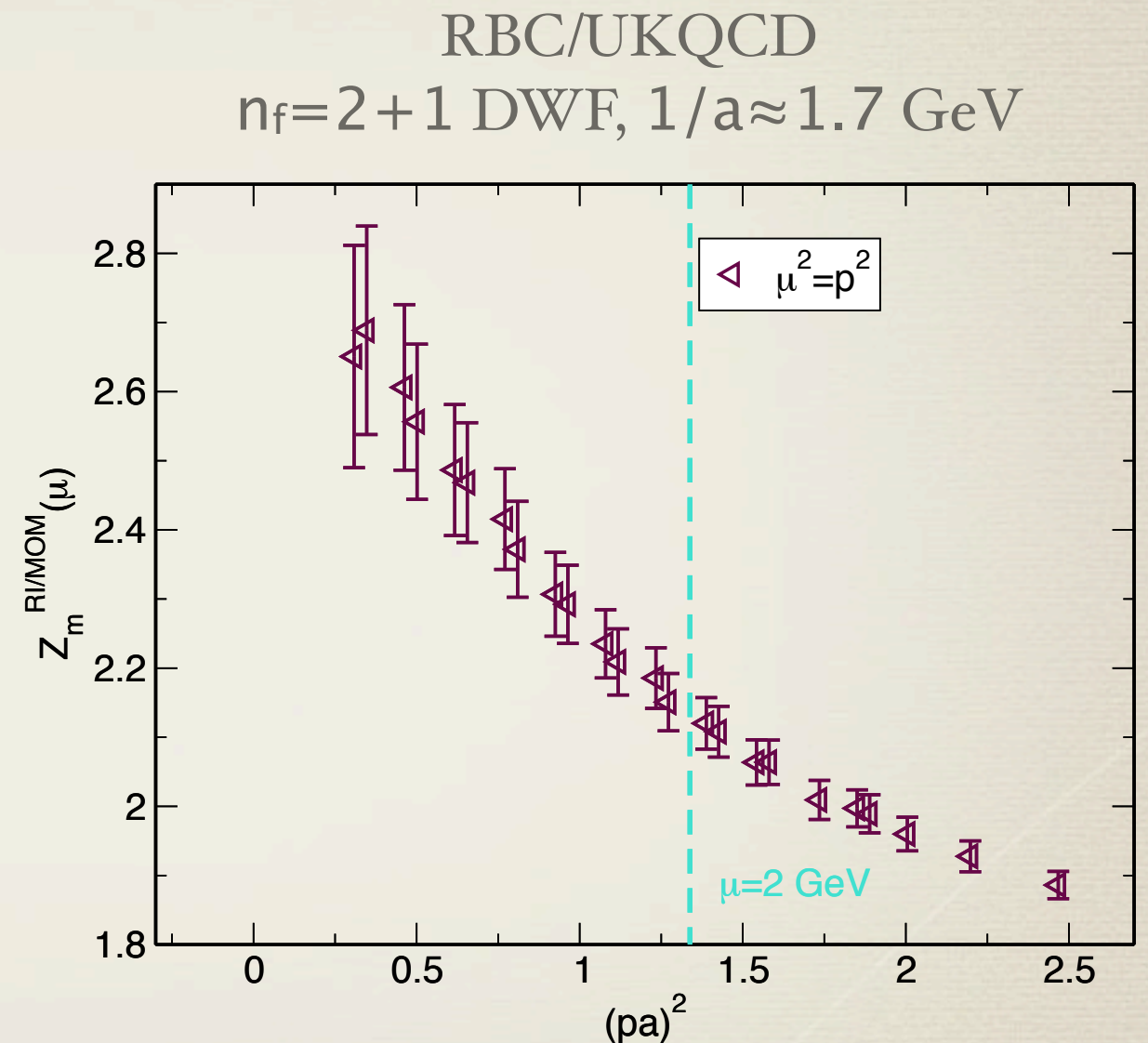


# RI/MOM procedure

- \*  $Z_m^{\text{RI}}(\mu)$  non-perturbative  $\mu$  dependence
- \*  $Z_m^{\text{RI}}(2 \text{ GeV})$  from  $p=2 \text{ GeV}$  intercept

$$Z_m^{\overline{MS}}(\mu) = \left( \frac{Z_m^{\overline{MS}}(\mu)}{Z_m^{\text{RI}}(\mu)} \right)_{PT} Z_m^{\text{RI,NPR}}(\mu)$$

- \* large momentum: discretization error?





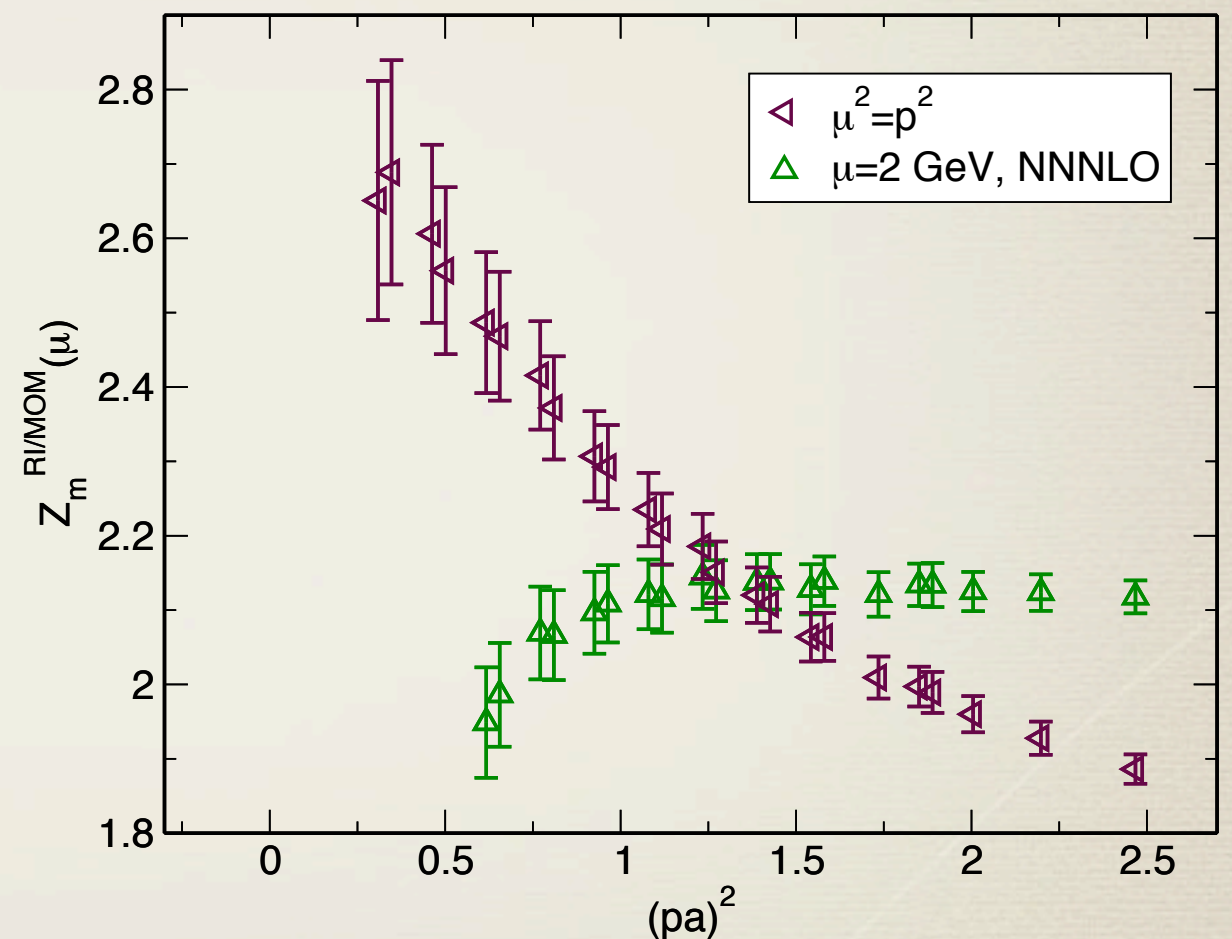
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- \* use PT to remove the running

RBC/UKQCD  
 $n_f=2+1$  DWF,  $1/a \approx 1.7 \text{ GeV}$



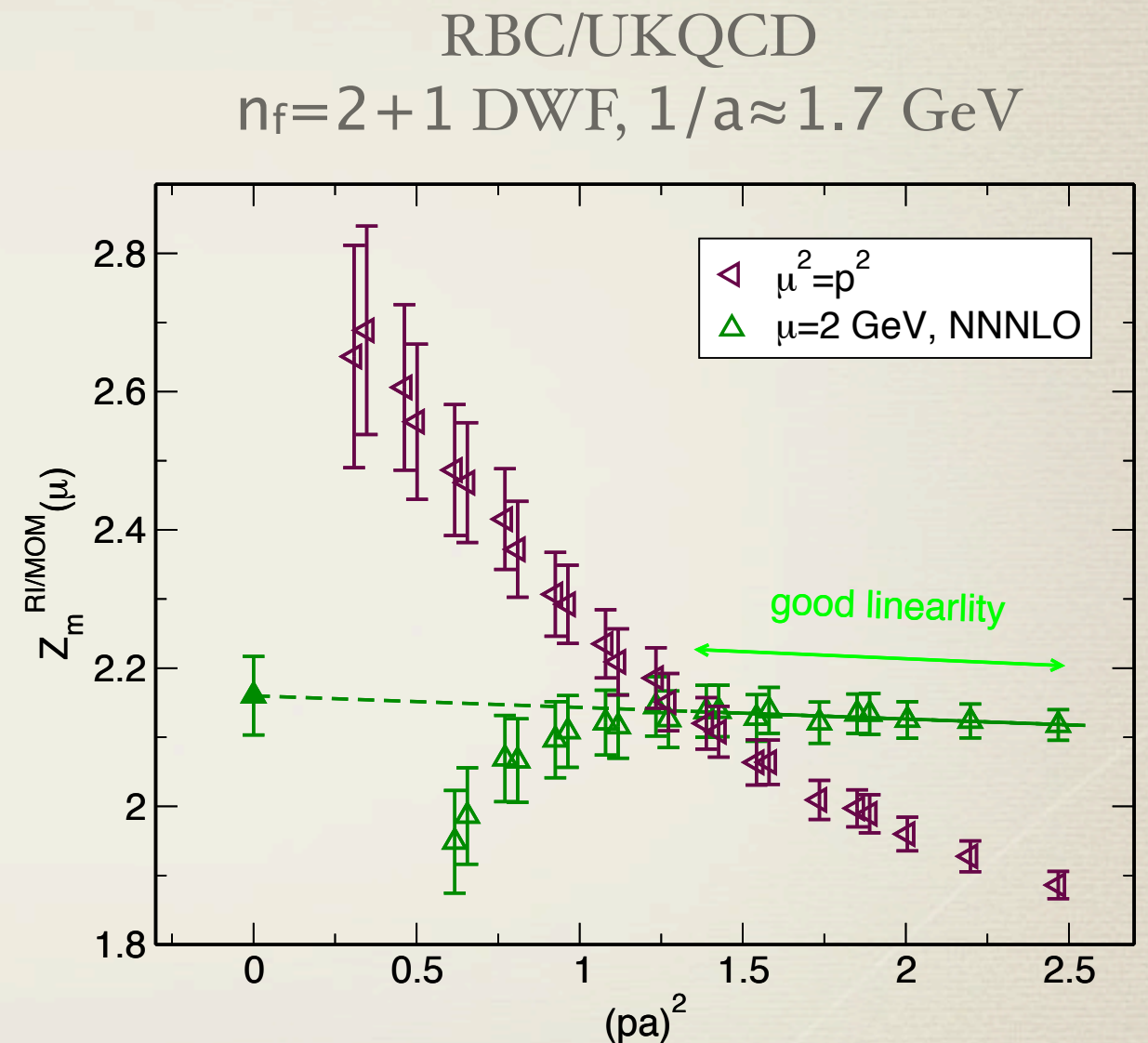


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- \*  $(pa)^2 \rightarrow 0$  using data in the window



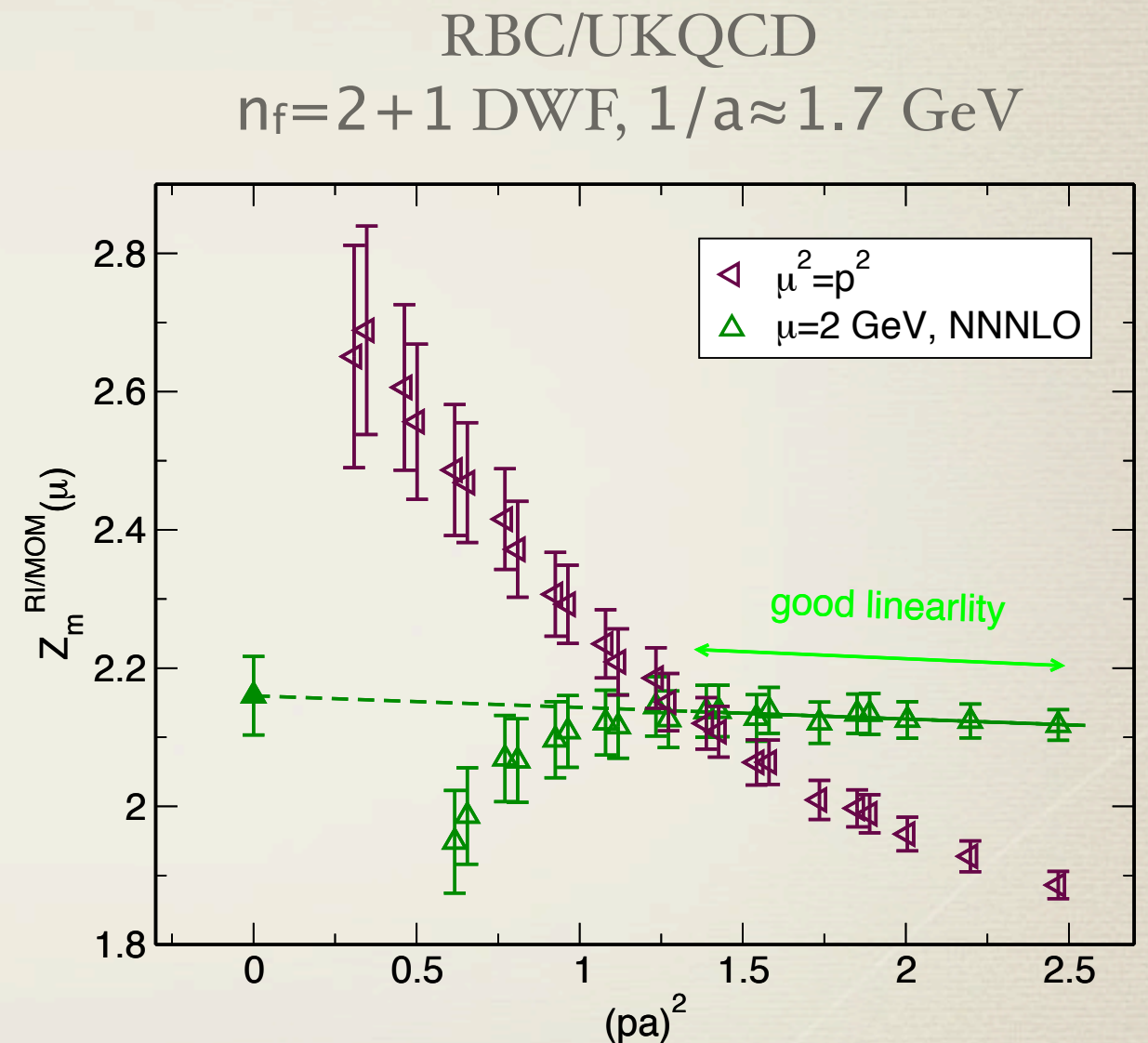


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- \* What if only known to NLO ?  
(before Franco and Lubicz 1998)



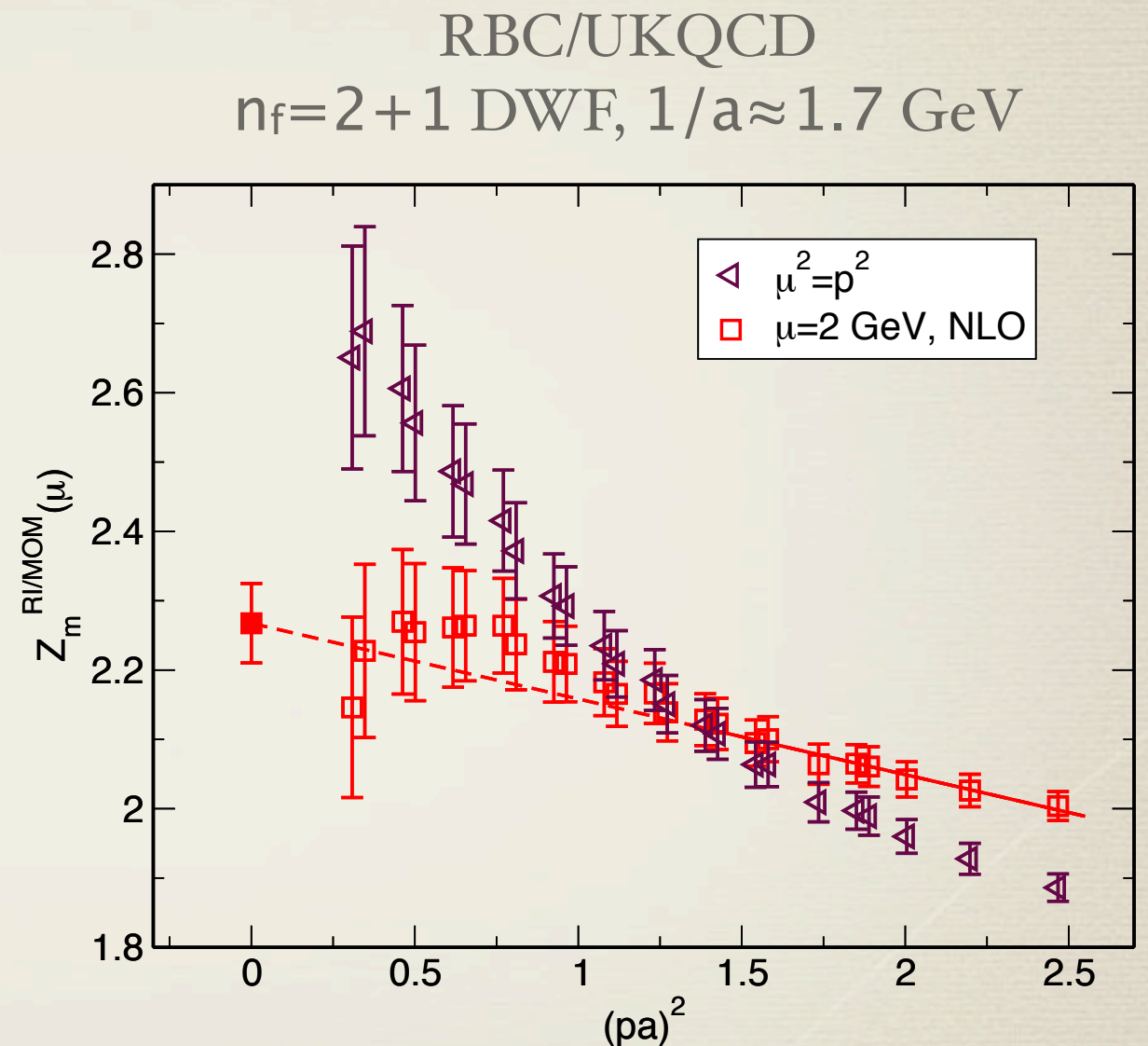


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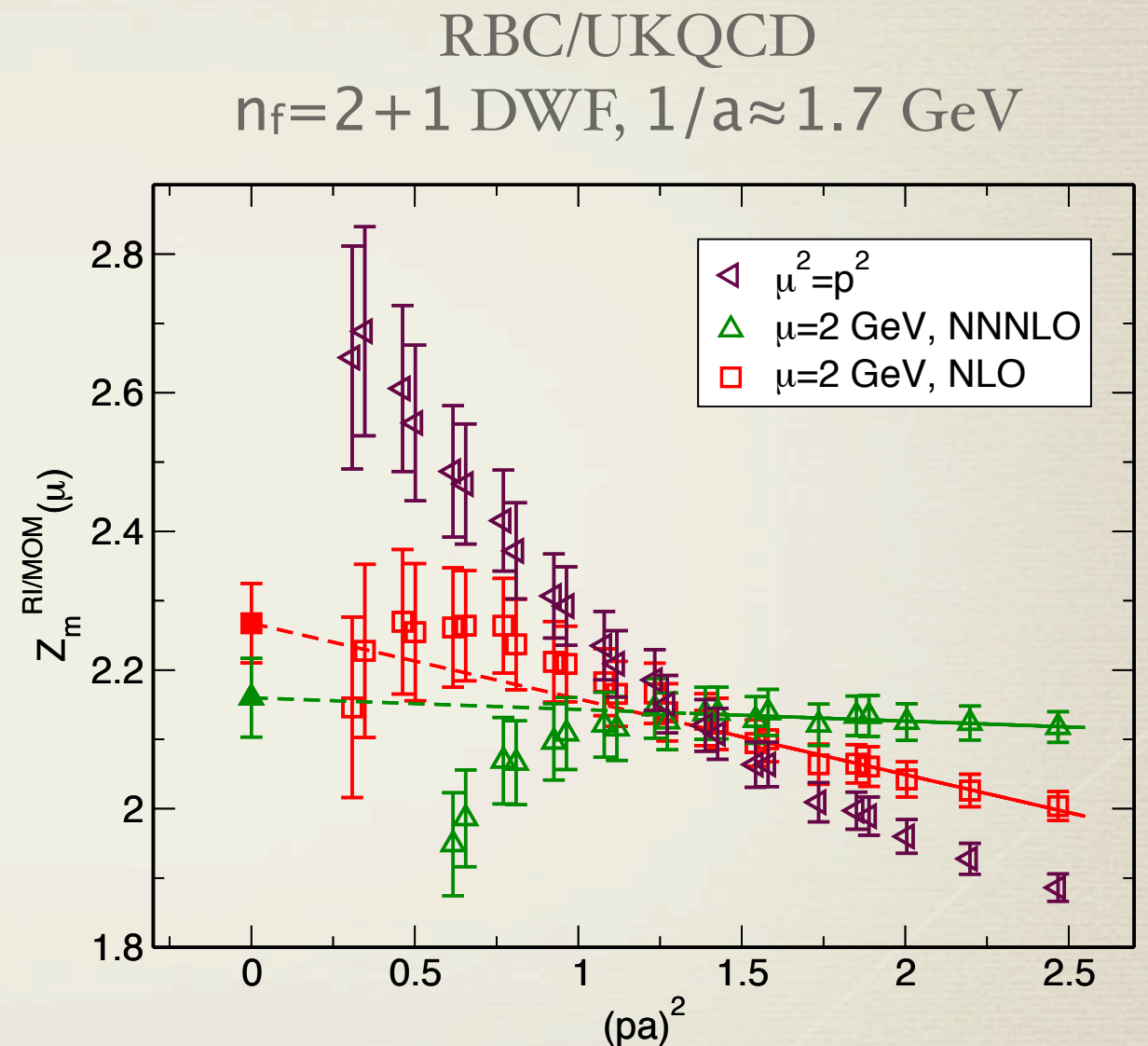


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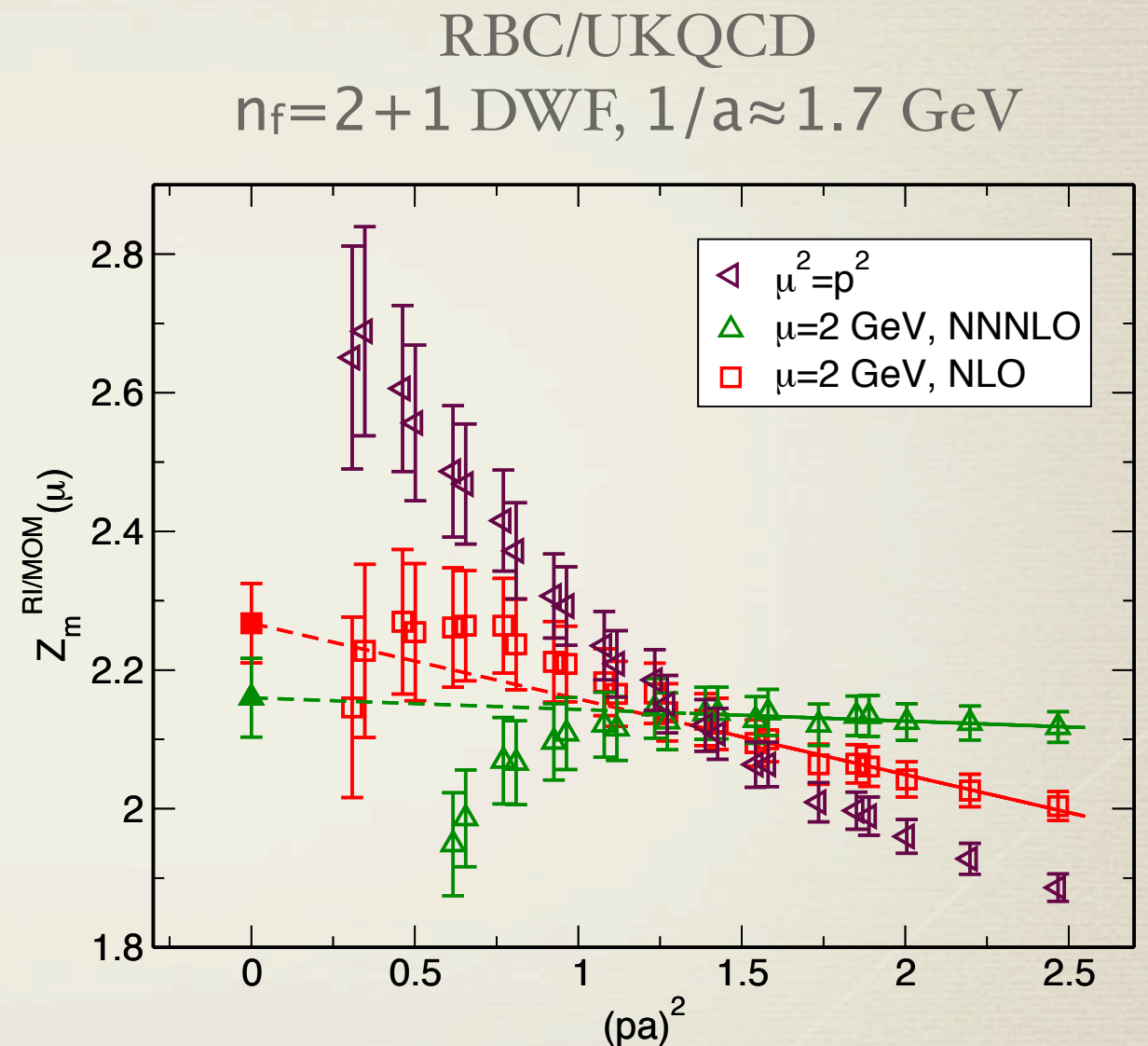


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  - linear extrapolation wrong!
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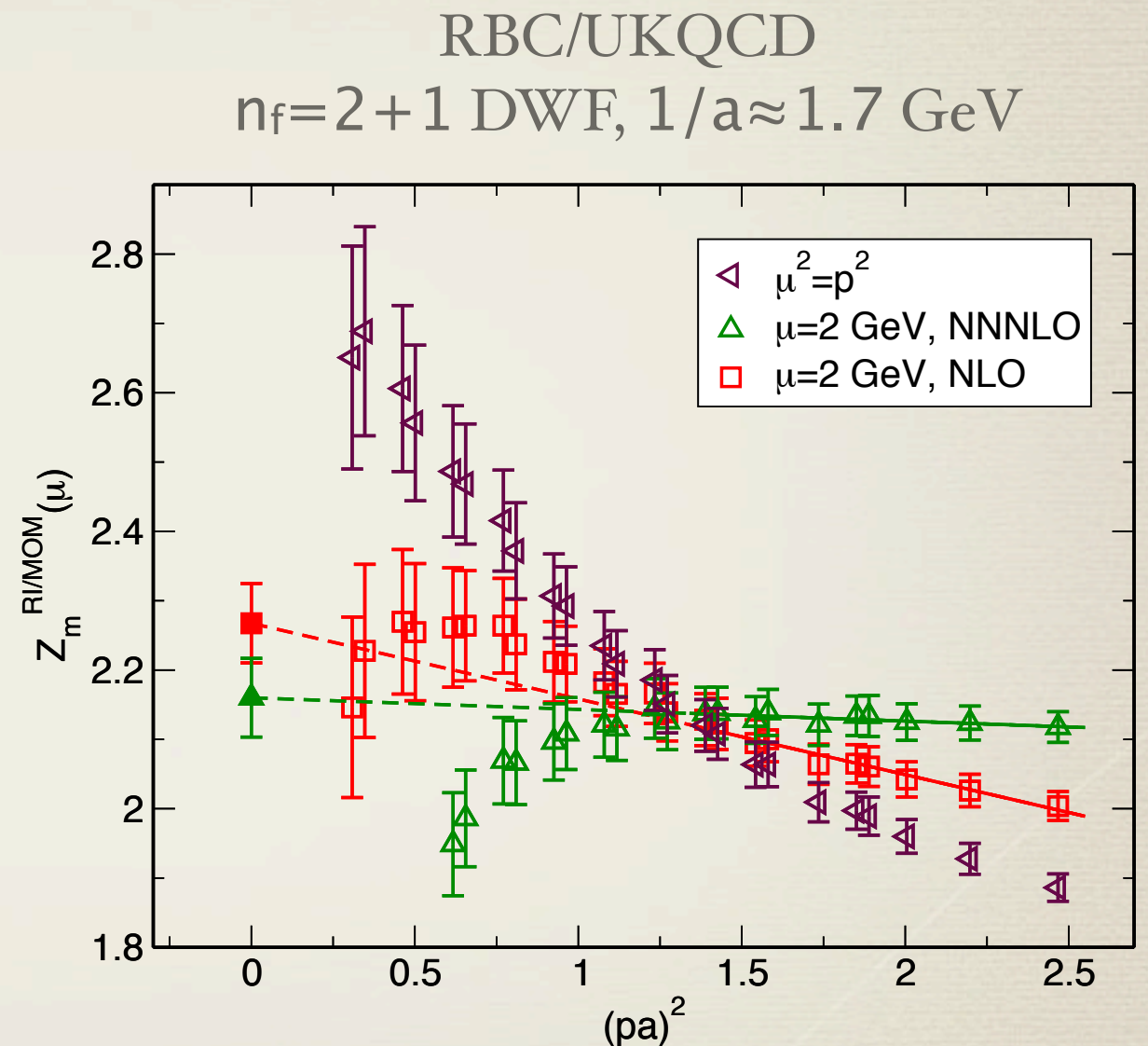
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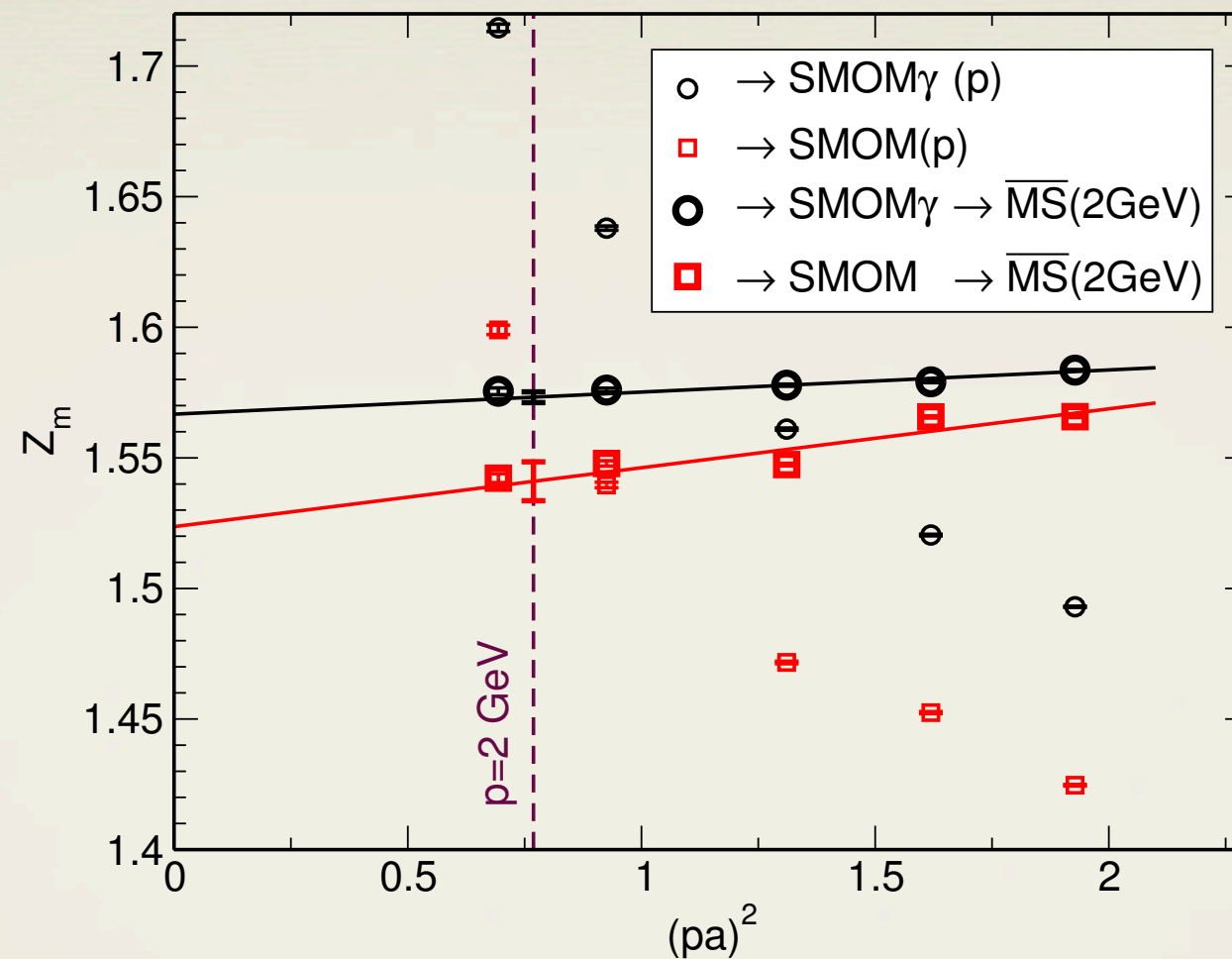
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- \* What if only known to NLO ?  
(before Franco and Lubicz 1998)

- linear extrapolation wrong!
- Be careful if you observe curvature or large slope
- a way out: no extrapolation, take  $p=2 \text{ GeV}$  value & add variation  $2\text{GeV} \rightarrow 0$  to systematic error





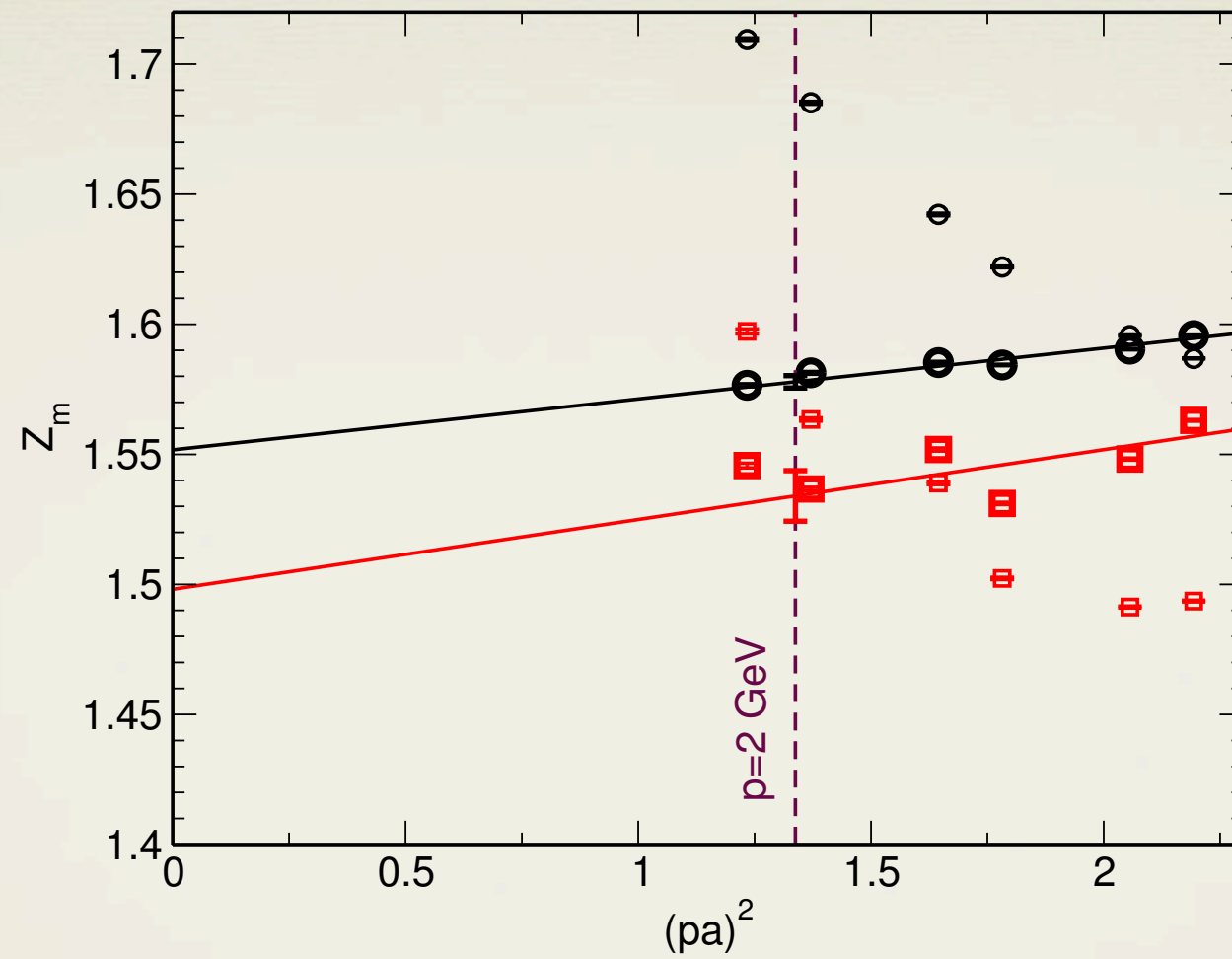
# SMOM mass renormalization: $32^3$



- \* larger  $O(4)$  breaking observed for SMOM
- \* take  $\text{SMOM}_\gamma$  for the intermediate scheme, use SMOM for sys error estimate



# SMOM mass renormalization: $24^3$





# SMOM results

$$Z_m^{\overline{\text{MS}}(32)}(\mu = 2 \text{ GeV}, n_f = 3; \text{SMOM}_{\gamma_\mu}) = 1.573(2),$$

$$Z_m^{\overline{\text{MS}}(24)}(\mu = 2 \text{ GeV}, n_f = 3; \text{SMOM}_{\gamma_\mu}) = 1.578(2),$$

- \* taking the **continuum limit** of  $Z(32)$ ,  $Z(24)/Z_1$  or  $Z(24)/Z_h$   
(to get rid of  $O(p^2 a^2)$  error)

$$Z_{mh}^{\overline{\text{MS}}(32)^c}(\mu = 2 \text{ GeV}, n_f = 3; \text{SMOM}_{\gamma_\mu}) = 1.510(6),$$

$$Z_{mh}^{\overline{\text{MS}}(32)^c}(\mu = 2 \text{ GeV}, n_f = 3; \text{SMOM}) = 1.495(22)$$

- \* sys error due to PT truncation may be estimated from the difference or the size of the highest order (2-loop) term of PT matching



# $Z_m$ error budget

ensemble	fine ( $32^2$ )	coarse ( $24^3$ )	coarse ( $16^3$ )[13]
intermediate scheme	RI/SMOM	RI/SMOM	RI/MOM
PT truncation error	2.1%	2.1%	6%
$m_s \neq 0$	0.1%	0.2%	7%
$(\Lambda_P - \Lambda_S)/2$	0.5%	0.6%	N.A. ( $\infty$ )
$(\Lambda_A - \Lambda_V)/2$	0.0%	0.0%	1%
total	2.2%	2.2%	9%



# the light quark mass results

$$\begin{aligned} m_{ud}^{\overline{\text{MS}}}(2\text{GeV}) &= Z_{ml}^{\overline{\text{MS}}(32)^c}(\mu = 2\text{GeV}, n_f = 3) \cdot \tilde{m}_{ud}(32^3) \cdot a^{-1}(32^3) \\ &= 3.59(13)_{\text{stat}}(14)_{\text{sys}}(8)_{\text{ren}} \text{ MeV}, \end{aligned}$$

$$\begin{aligned} m_s^{\overline{\text{MS}}}(2\text{GeV}) &= Z_{mh}^{\overline{\text{MS}}(32)^c}(\mu = 2\text{GeV}, n_f = 3) \cdot \tilde{m}_s(32^3) \cdot a^{-1}(32^3) \\ &= 96.2(1.6)_{\text{stat}}(0.2)_{\text{sys}}(2.1)_{\text{ren}} \text{ MeV}, \end{aligned}$$



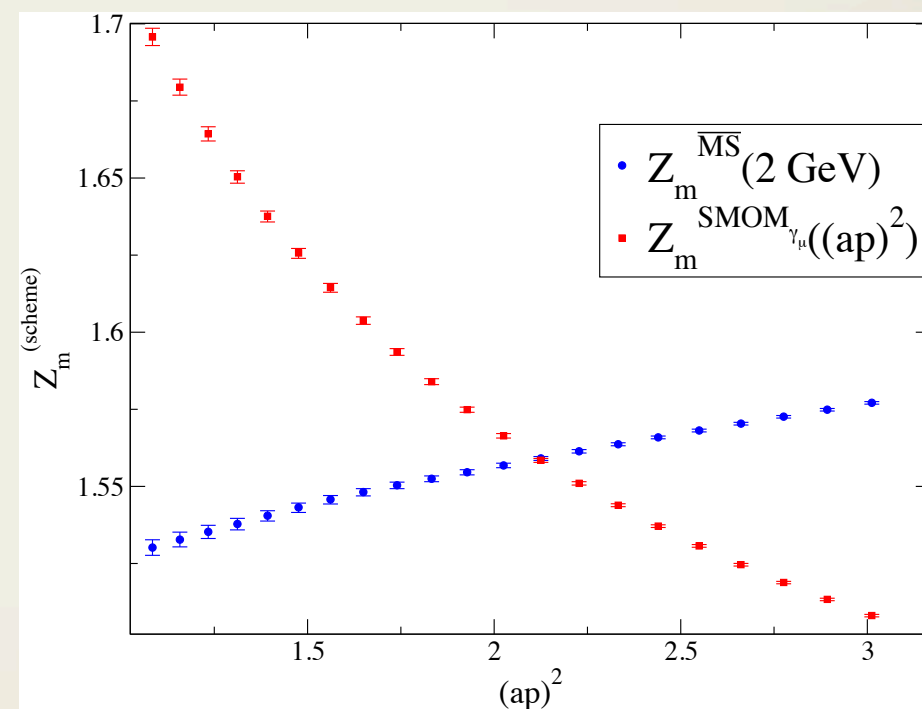
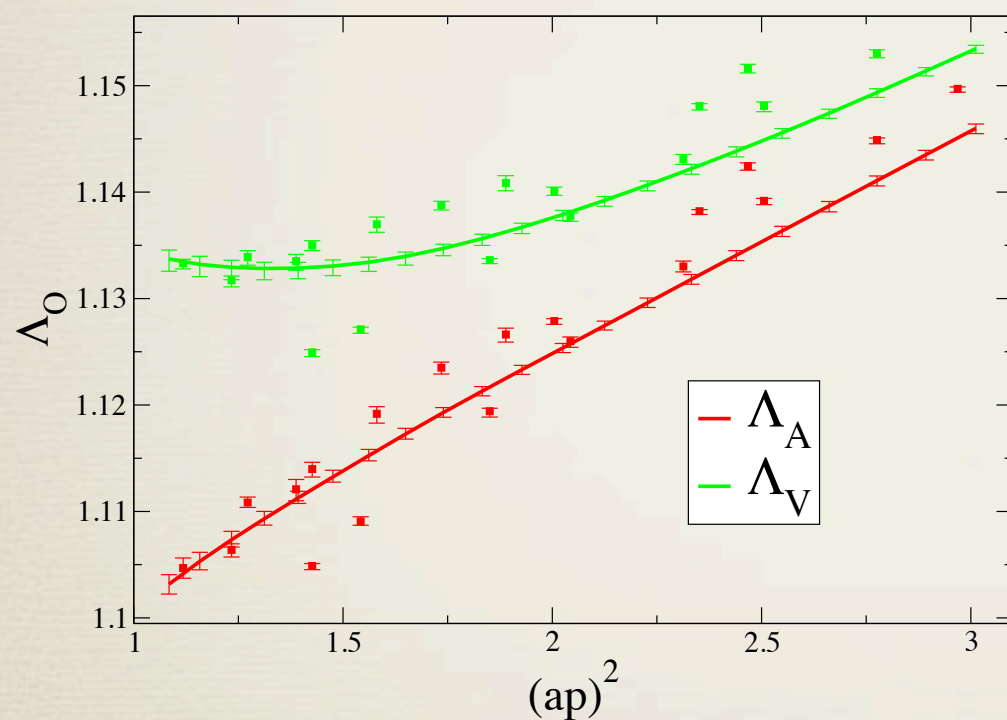
# Outlook

- \* more chiral DSDR:
  - \* smaller masses (down to unitary  $m_\pi \approx 180$  MeV) larger volume, but, coarser lattice
  - \* with twisted mass to reduce  $m_{\text{res}}$
- \* in principle, just adding another ensemble with different  $a^2$
- \* no need of New NPR if  $32^3$  is used again as the primary ensemble
- \* but...



# Outlook

- \* promising new idea of twisted boundary RI/MOM scheme by R. Arthur, P. Boyle.
- \* fix base momentum & twist boundary to smoothly change the momentum
- ➔ no  $O(4)$  breaking effect in  $p^2$  dependence → much precise handling possible



- \* step scaling tested
- \* Further improvements are expected!!!



VIELEN DANK!